# ANISOTROPIC $(\vec{p}, \vec{q})$-LAPLACIAN PROBLEMS WITH SUPERLINEAR NONLINEARITIES 

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#### Abstract

In this paper we consider a class of anisotropic ( $\vec{p}, \vec{q}$ )-Laplacian problems with nonlinear right-hand sides that are superlinear at $\pm \infty$. We prove the existence of two nontrivial weak solutions to this kind of problem by applying an abstract critical point theorem under very general assumptions on the data without supposing the Ambrosetti-Rabinowitz condition.


## 1. Introduction

Let $\Omega \subseteq \mathbb{R}^{N}, N \geq 2$ be a bounded domain with boundary $\partial \Omega$ of class $C^{1}$. In this paper we consider the following anisotropic differential equation involving the $(\vec{p}, \vec{q})$-Laplacian given by

$$
\begin{aligned}
-\Delta_{\vec{p}} u-\Delta_{\vec{q}} u & =\lambda f(x, u) & & \text { in } \Omega, \\
u & =0 & & \text { on } \partial \Omega,
\end{aligned}
$$

where $\lambda$ is a positive parameter, $\vec{p}=\left(p_{1}, p_{2}, \ldots, p_{N}\right), \vec{q}=\left(q_{1}, q_{2}, \ldots, q_{N}\right), \vec{p}, \vec{q} \in \mathbb{R}^{N}$ are real vectors such that

$$
\max \left\{q_{1}, q_{2}, \ldots, q_{N}\right\}<\min \left\{p_{1}, p_{2}, \ldots, p_{N}\right\}<N
$$

and $f: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a nonlinearity with subcritical growth that is superlinear at $\pm \infty$, see $\left(\mathrm{H}_{\mathrm{f}}\right)$ for the precise assumptions. Moreover, for any $\vec{s}=\left(s_{1}, s_{2}, \ldots, s_{N}\right) \in$ $\mathbb{R}^{N}$, we denote by

$$
\Delta_{\vec{s}} u=\sum_{i=1}^{N} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{s_{i}-2} \frac{\partial u}{\partial x_{i}}\right)
$$

the anisotropic $\vec{s}$-Laplace differential operator. If $s_{i}=2$ for all $i=1, \ldots, N$, we get

$$
\sum_{i=1}^{N} \frac{\partial^{2} u}{\partial x_{i}^{2}}=\Delta u
$$

that is the usual Laplacian and if $\vec{s}$ is constant (i.e. $s_{i}=s$ for each $i=1, \ldots, N$ ) we obtain

$$
\sum_{i=1}^{N} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{s-2} \frac{\partial u}{\partial x_{i}}\right)=\tilde{\Delta}_{s} u
$$

which is called the pseudo-s-Laplace operator, see Belloni-Kawohl [3] or BrascoFranzina [10]. More information about anisotropic operators and in particular

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about the theory of anisotropic Sobolev spaces can be found in Kufner-Rákosník [20], Nikol'skiŭ [24] and Rákosník [29, 30].

Anisotropic differential problems have a large background in several applications, for example, the study of an epidemic disease in heterogeneous habitat can be expressed by an anisotropic nonlinear system. In general, anisotropic operators are used for modeling in which partial differential derivatives vary with the direction. Also, the anisotropic Laplacian, given by

$$
\sum_{i=1}^{N} \frac{\partial}{\partial x_{i}}\left(\frac{\partial}{\partial \xi_{i}}\left(\frac{1}{2} F^{2}\right)(\nabla u)\right)
$$

where $F(\xi)=\left(\sum_{i=1}^{N}\left|\xi_{i}\right|^{2}\right)^{\frac{1}{2}}$ for $\xi \in \mathbb{R}^{N}$, plays a key role in physical models of crystal growth in the context of the so-called Wulff shape of $F$ (also known as equilibrium crystal shape), see the work of Wulff [39]. For more information on applications in different disciplines we refer to Antontsev-Díaz-Shmarev [1], Bendahmane-Chrif-El Manouni [4], Bendahmane-Langlais-Saad [5], Vétois [37] and the references therein.

Although there are some works for $\vec{p}$-Laplacian problems, only a few exist for the anisotropic $(\vec{p}, \vec{q})$-Laplacian. Recently, Razani-Figueiredo [32] studied the anisotropic Dirichlet problem

$$
\begin{equation*}
-\sum_{i=1}^{N} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial u}{\partial x_{i}}\right)-\sum_{i=1}^{N} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial u}{\partial x_{i}}\right)=\lambda u^{\gamma-1} \tag{1.1}
\end{equation*}
$$

in a bounded and regular domain $\Omega$ of $\mathbb{R}^{N}$ with $\gamma>1$ and $\lambda>0$. Based on a sub-supersolution approach the authors show the existence of at least one positive solution of (1.1). In Razani [31] nonstandard competing anisotropic ( $\vec{p}, \vec{q}$ )-Laplacian problems with convolution of the form

$$
\begin{equation*}
-\sum_{i=1}^{N} \frac{\partial}{\partial x_{i}}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}-2}-\mu\left|\frac{\partial u}{\partial x_{i}}\right|^{q_{i}-2}\right) \frac{\partial u}{\partial x_{i}}=f(x, \phi * u, \nabla(\phi * u)) \tag{1.2}
\end{equation*}
$$

have been considered, where $\phi \in L^{1}\left(\mathbb{R}^{N}\right)$. If $\mu>0$, then the existence of a generalized solution of (1.2) is shown by using a Galerkin base for the function space and if $\mu \leq 0$, then any generalized solution turns out to be a weak solution. We also mention the recent work of Tavares [36] who considered existence and multiplicity of nonnegative solutions for the problem given by

$$
\begin{aligned}
-\Delta_{\vec{p}} u-\Delta_{\vec{q}} u & =k(x) u^{\alpha-1}+f(x, u) & & \text { in } \Omega, \\
u & =0 & & \text { on } \partial \Omega
\end{aligned}
$$

where $\alpha>1, k \in L^{\infty}(\Omega)$ with $k(x)>0$ for a.a. $x \in \Omega$ and a Carathéodory nonlinearity $f: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ with subcritical growth. Existence and regularity results for anisotropic problems driven by the $\vec{p}$-Laplacian have been obtained by several authors. Without guarantee of completeness, we mention just some of them and refer to the papers of Bonanno-D'Aguì-Sciammetta [9], Ciani-FigueiredoSuárez [11], Ciraolo-Figalli-Roncoroni [12], Ciraolo-Sciammetta [13], DiBenedetto-Gianazza-Vespri [14], dos Santos-Figueiredo-Tavares [15], Fragalà-Gazzola-Kawohl [17], Perera-Agarwal-O'Regan [27], Ragusa-Razani-Safari [28], see also the references therein. Related results for the $(p, q)$-Laplacian, double phase equations, anisotropic problems or the discrete $p$-Laplacian can be found in the works of

Bai-Papageorgiou-Zeng [2], Bohner-Caristi-Ghobadi-Heidarkhani [6], El Manouni-Marino-Winkert [16], Leggat-Miri [21], He-Ousbika-El Allali-Zuo [18], Ju-Molica Bisci-Zhang [19], Liu-Motreanu-Zeng [23], Papageorgiou [25], Son-Sim [34], VetroVetro [38], Zeng-Bai-Gasiński-Winkert [40], Zeng-Papageorgiou [41] and Zeng-Rădu-lescu-Winkert [42].

Motivated by the above mentioned works for $\vec{p}$-Laplacian problems, we consider in our paper so-called $(\vec{p}, \vec{q})$-Laplacian problems with general right-hand sides. Our main goal is to apply a two critical point result due to Bonanno-D'Aguì [7, Theorem 2.1] in order to get the existence of two positive solutions for problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ with different energy sign. Note that the results in [7] are given in a very general setting which can be applied to a large number of problems. Our paper can be seen as an extension of the work of Bonanno-D'Aguì-Sciammetta [8] for ( $\vec{p}, \vec{q}$ )-Laplacian problems. But not only the differential operator is more general than in [8], also the condition required on $f$ are weaker. Indeed, instead of assuming the AmbrosettiRabinowitz condition, we only assume that the nonlinear term on the right-hand side of $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ is $\left(p^{+}-1\right)$-superlinear at $\pm \infty$ (with $p^{+}$defined in (2.2) replacing $h$ by $p$ ) and fulfills in addition a suitable behavior at $\pm \infty$, see $\left(\mathrm{H}_{\mathrm{f}}\right)$. These hypotheses are weaker than the Ambrosetti-Rabinowitz condition. Under these conditions, together with the subcritical growth, we prove the existence of two weak solutions for problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ that have opposite energy sign related to the energy functional of $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$.

The paper is organized as follows. In Section 2 we present the main properties of anisotropic Sobolev spaces and consider the main features of the anisotropic $(\vec{p}, \vec{q})$ Laplacian, see Propositions 2.3 and 2.4. Moreover, we recall the main abstract critical point theorem which will play a key role in our treatment, see Theorem 2.7. In Section 3 we first state the precise assumptions on the data of problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ and prove that the corresponding energy functional fulfills the C-condition. Then we give our main result about the existence of two nontrivial weak solutions of $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ (see Theorem 3.3) and some consequences in which the solutions are nonnegative, see Theorems 3.4 and 3.5. Finally, we consider the autonomous problem $\left(A D_{\lambda}^{\vec{p}, \vec{q}}\right)$, providing an existence result (see Corollary 3.6) and an example.

## 2. Preliminaries and basic properties

Let $\Omega \subset \mathbb{R}^{N}, N \geq 2$, be a bounded domain with boundary $\partial \Omega$ of class $C^{1}$. For any real vector $\vec{h}=\left(h_{0}, h_{1}, h_{2}, \ldots, h_{N}\right)$ with $h_{i} \geq 1$ for every $i=0,1, \ldots, N$, we indicate with $W^{1, \vec{h}}(\Omega)$ the anisotropic Sobolev space defined by

$$
W^{1, \vec{h}}(\Omega)=\left\{u \in L^{h_{0}}(\Omega): \frac{\partial u}{\partial x_{i}} \in L^{h_{i}}(\Omega) \text { for } i=1, \ldots, N\right\}
$$

equipped with the norm

$$
\begin{equation*}
\|u\|_{W^{1, \vec{h}}(\Omega)}=\|u\|_{L^{h_{0}}(\Omega)}+\sum_{i=1}^{N}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{h_{i}}(\Omega)} \tag{2.1}
\end{equation*}
$$

Moreover, set

$$
\begin{equation*}
h^{-}=\min \left\{h_{1}, h_{2}, \ldots, h_{N}\right\} \quad \text { and } \quad h^{+}=\max \left\{h_{1}, h_{2}, \ldots, h_{N}\right\}, \tag{2.2}
\end{equation*}
$$

and suppose that $h^{-}<N$ and $h_{0}<\left(h^{-}\right)^{*}=\frac{N h^{-}}{N-h^{-}}$. Denote by $W_{0}^{1, \vec{h}}(\Omega)$ the closure of $C_{0}^{\infty}(\Omega)$ endowed with the following norm

$$
\begin{equation*}
\|u\|_{W_{0}^{1, \vec{h}}(\Omega)}:=\sum_{i=1}^{N}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{h_{i}}(\Omega)} \tag{2.3}
\end{equation*}
$$

which is equivalent to the usual one given in (2.1). Indeed, taking into account that $W^{1, h^{-}}(\Omega)$ is compactly embedded in $L^{h_{0}}(\Omega)$ and using Hölder's inequality (see (2.4) in Proposition 2.1), we have that

$$
\|u\|_{L^{h_{0}}(\Omega)} \leq c\|u\|_{W^{1, h^{-}}(\Omega)} \leq \tilde{k} \sum_{i=1}^{N}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{h_{i}}(\Omega)}
$$

It is well known that $W_{0}^{1, \vec{h}}(\Omega)$, endowed with the norm defined in (2.3), is a separable Banach space and it is also reflexive if $h_{i}>1$ for all $i=1, \ldots, N$, see Rákosník [29, Theorem 1].

Given $\vec{p}=\left(p_{0}, p_{1}, p_{2}, \ldots, p_{N}\right)$ and $\vec{q}=\left(q_{0}, q_{1}, q_{2}, \ldots, q_{N}\right)$, with $p_{i}, q_{i} \geq 2$ for every $i=1, \ldots, N$, we suppose that
(H) $q^{+}<p^{-}<N, p_{0}<\left(p^{-}\right)^{*}$ and $q_{0}<\left(q^{-}\right)^{*}$, where $(\cdot)^{*}=\frac{N(\cdot)}{N-(\cdot)}$ denotes the critical Sobolev exponent.
In the following proposition, we give a relation between the spaces $W_{0}^{1, \vec{p}}(\Omega)$ and $W_{0}^{1, \vec{q}}(\Omega)$ and their norms. In particular, we underline that $p_{0}$ and $q_{0}$ are necessary only for the definition of the anisotropic Sobolev spaces, but since we endow the spaces with the equivalent norm given in (2.3), from now on we will only deal with the components $\left(p_{1}, \ldots, p_{N}\right)$ and $\left(q_{1}, \ldots, q_{N}\right)$.
Proposition 2.1. If $q^{+}<p^{-}$, then $W_{0}^{1, \vec{p}}(\Omega) \subseteq W_{0}^{1, \vec{q}}(\Omega)$ and

$$
\|u\|_{W_{0}^{1, \vec{q}}(\Omega)} \leq \max _{1 \leq i \leq N}\left\{|\Omega|^{\frac{p_{i}-q_{i}}{p_{i} q_{i}}}\right\}\|u\|_{W_{0}^{1, \vec{p}}(\Omega)}
$$

where $|\Omega|$ is the Lebesgue measure of $\Omega$.
Proof. Fix $u \in W_{0}^{1, \vec{p}}(\Omega)$ and $i \in\{1, \ldots, N\}$. In particular, $\frac{\partial u}{\partial x_{i}} \in L^{p_{i}}(\Omega)$ and $\left|\frac{\partial u}{\partial x_{i}}\right|^{p^{-}} \in L^{\frac{p_{i}}{p^{-}}}(\Omega)$. If $p_{i}>p^{-}$, by Hölder's inequality, we get

$$
\begin{equation*}
\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p^{-}} \mathrm{d} x\right)^{\frac{1}{p^{-}}} \leq|\Omega|^{\frac{p_{i}-p^{-}}{p_{i} p^{-}}}\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x\right)^{\frac{1}{p_{i}}} \tag{2.4}
\end{equation*}
$$

while if $p_{i}=p^{-}$, then (2.4) is an equality. Moreover, $\left|\frac{\partial u}{\partial x_{i}}\right|^{q^{+}} \in L^{\frac{p^{-}}{q^{+}}}(\Omega)$, then again from Hölder's inequality, we obtain

$$
\begin{equation*}
\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{q^{+}} \mathrm{d} x\right)^{\frac{1}{q^{+}}} \leq|\Omega|^{\frac{p^{-}-q^{+}}{p^{-} q^{+}}}\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p^{-}} \mathrm{d} x\right)^{\frac{1}{p^{-}}} \tag{2.5}
\end{equation*}
$$

Thus, combining (2.4) and (2.5) we derive

$$
\begin{equation*}
\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{q^{+}} \mathrm{d} x\right)^{\frac{1}{q^{+}}} \leq|\Omega|^{\frac{p_{i}-q^{+}}{p_{i} q^{+}}}\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x\right)^{\frac{1}{p_{i}}} \tag{2.6}
\end{equation*}
$$

for all $i=1, \ldots, N$. Furthermore, if $q^{+}>q_{i}$ and using Hölder's inequality we have

$$
\begin{equation*}
\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{q_{i}} \mathrm{~d} x\right)^{\frac{1}{q_{i}}} \leq|\Omega|^{\frac{q^{+}-q_{i}}{q^{+} q_{i}}}\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{q^{+}} \mathrm{d} x\right)^{\frac{1}{q^{+}}} \tag{2.7}
\end{equation*}
$$

and if $q^{+}=q_{i}$ the previous inequality becomes an equality. Then, from (2.6) and (2.7) it follows that

$$
\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{q_{i}} \mathrm{~d} x\right)^{\frac{1}{q_{i}}} \leq|\Omega|^{\frac{p_{i}-q_{i}}{p_{i} q_{i}}}\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x\right)^{\frac{1}{p_{i}}} \quad \text { for all } i=1, \ldots, N
$$

Hence,

$$
\|u\|_{W_{0}^{1, \vec{q}}(\Omega)} \leq \max _{1 \leq i \leq N}\left\{|\Omega|^{\frac{p_{i}-q_{i}}{p_{i} q_{i}}}\right\}\|u\|_{W_{0}^{1, \vec{p}}(\Omega)}
$$

and the proof is complete.
In the sequel, we estimate the embedding constant of $W_{0}^{1, \vec{p}}(\Omega)$ into $L^{r}(\Omega)$ for each $r \in\left[1,\left(p^{-}\right)^{*}\right]$ with $p^{-}<N$, where $\left(p^{-}\right)^{*}$ is the critical Sobolev exponent to $p^{-}$, that is

$$
\begin{equation*}
\left(p^{-}\right)^{*}=\frac{N p^{-}}{N-p^{-}} . \tag{2.8}
\end{equation*}
$$

Proposition 2.2. If $1 \leq p^{-}<N$, then for any $r \in\left[1,\left(p^{-}\right)^{*}\right]$, $W_{0}^{1, \vec{p}}(\Omega) \hookrightarrow L^{r}(\Omega)$ is continuous and one has

$$
\begin{equation*}
\|u\|_{L^{r}(\Omega)} \leq T_{r}\|u\|_{W_{0}^{1, \vec{p}}(\Omega)} \tag{2.9}
\end{equation*}
$$

for all $u \in W_{0}^{1, \vec{p}}(\Omega)$, where

$$
\begin{align*}
& T_{r}=c_{r} 2^{\frac{(N-1)\left(p^{-}-1\right)}{p^{-}}} \max _{1 \leq i \leq N}\left\{|\Omega|^{\frac{p_{i}-p^{-}}{p_{i} p^{-}}}\right\} \\
& c_{r}=T|\Omega|^{\frac{\left(p^{-}\right)^{*-r}}{\left(p^{-}\right)^{* r}}},  \tag{2.10}\\
& T \leq \frac{N^{-\frac{1}{p^{-}}}}{\sqrt{\pi}}\left(\frac{p^{-}-1}{N-p^{-}}\right)^{1-\frac{1}{p^{-}}}\left(\frac{\Gamma(N) \Gamma\left(1+\frac{N}{2}\right)}{\Gamma\left(\frac{N}{p^{-}}\right) \Gamma\left(1+N-\frac{N}{p^{-}}\right)}\right)^{\frac{1}{N}}
\end{align*}
$$

see Talenti [35, formula (2)] and $\Gamma$ is the Euler function. Moreover, for any $r \in$ $\left[1,\left(p^{-}\right)^{*}\left[\right.\right.$ the embedding $W_{0}^{1, \vec{p}}(\Omega) \hookrightarrow L^{r}(\Omega)$ is compact.
Proof. From the Sobolev embedding theorem there exists a positive constant $T \in \mathbb{R}$ such that for all $u \in W_{0}^{1, p^{-}}(\Omega)$ the following holds

$$
\begin{equation*}
\|u\|_{L^{\left(p^{-}\right)^{*}}(\Omega)} \leq T\|u\|_{W_{0}^{1, p^{-}}(\Omega)}, \tag{2.11}
\end{equation*}
$$

where $\left(p^{-}\right)^{*}$ and $T$ are defined in (2.8) and (2.10), respectively.
Fix $r \in\left[1,\left(p^{-}\right)^{*}\right]$. Since $\frac{1}{r}=\frac{1}{\left(p^{-}\right)^{*}}+\frac{\left(p^{-}\right)^{*}-r}{\left(p^{-}\right)^{*} r}$, by Hölder's inequality and (2.11), we have

$$
\|u\|_{L^{r}(\Omega)} \leq\|u\|_{L^{\left(p^{-}\right)^{*}}(\Omega)}|\Omega|^{\frac{\left(p^{-}\right)^{*}-r}{\left(p^{-}\right)^{*} r}} \leq T|\Omega|^{\frac{\left(p^{-}\right)^{*}-r}{\left(p^{-}\right)^{*} r}}\|u\|_{W_{0}^{1, p^{-}}(\Omega)}
$$

that is

$$
\begin{equation*}
\|u\|_{L^{r}(\Omega)} \leq c_{r}\|u\|_{W_{0}^{1, p^{-}}(\Omega)} . \tag{2.12}
\end{equation*}
$$

Arguing as in the proof of Proposition 2.1 in Bonanno-D'Aguì-Sciammetta [8], we derive

$$
\|u\|_{W_{0}^{1, p^{-}}(\Omega)} \leq 2^{\frac{(N-1)\left(p^{-}-1\right)}{p^{-}}} \max _{1 \leq i \leq N}\left\{\left\lvert\, \Omega^{\frac{p_{i}-p^{-}}{p_{i} p^{-}}}\right.\right\}\|u\|_{W_{0}^{1, \vec{p}}(\Omega)}
$$

and, taking (2.12) into account, we get that (2.9) holds for every $r \in\left[1,\left(p^{-}\right)^{*}\right]$.
Finally, combining the continuous embedding $W_{0}^{1, \vec{p}}(\Omega) \hookrightarrow W_{0}^{1, p^{-}}(\Omega)$ with the compact embedding $W_{0}^{1, p^{-}}(\Omega) \hookrightarrow L^{r}(\Omega)$, it follows that $W_{0}^{1, \vec{p}}(\Omega) \hookrightarrow L^{r}(\Omega)$ is compact for any $r \in\left[1,\left(p^{-}\right)^{*}[\right.$.

Now, we define

$$
F(x, t)=\int_{0}^{t} f(x, \xi) \mathrm{d} \xi \quad \text { for all }(x, t) \in \Omega \times \mathbb{R}
$$

and we introduce the functionals $\Phi, \Psi: W_{0}^{1, \vec{p}}(\Omega) \rightarrow \mathbb{R}$ given by

$$
\begin{align*}
& \Phi(u)=\sum_{i=1}^{N}\left(\frac{1}{p_{i}} \int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x+\frac{1}{q_{i}} \int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{q_{i}} \mathrm{~d} x\right)  \tag{2.13}\\
& \Psi(u)=\int_{\Omega} F(x, u) \mathrm{d} x
\end{align*}
$$

for every $u \in W_{0}^{1, \vec{p}}(\Omega)$. Clearly, $\Phi$ and $\Psi$ are Gâteaux differentiable and one has

$$
\begin{align*}
\Phi^{\prime}(u)(v) & =\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial u}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x+\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial u}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x  \tag{2.14}\\
\Psi^{\prime}(u)(v) & =\int_{\Omega} f(x, u) v \mathrm{~d} x
\end{align*}
$$

for every $u, v \in W_{0}^{1, \vec{p}}(\Omega)$. Also, we consider the energy functional $I_{\lambda}: W_{0}^{1, \vec{p}}(\Omega) \rightarrow \mathbb{R}$ associated to our problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$, that is given by $I_{\lambda}=\Phi-\lambda \Psi$ for all $\lambda>0$.

We recall that $u: \Omega \rightarrow \mathbb{R}$ is a weak solution of problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ if the following condition holds for all $v \in W_{0}^{1, \vec{p}}(\Omega)$

$$
\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial u}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x+\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial u}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x=\lambda \int_{\Omega} f(x, u) v \mathrm{~d} x
$$

Then, from (2.14) it follows that $u \in W_{0}^{1, \vec{p}}(\Omega)$ is a weak solution of problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ if and only if $u$ is a critical point for $I_{\lambda}$. Consequently, our study is based on critical point theory and in particular on a critical point theorem due to Bonanno-D'Aguì [7] that we state later in Theorem 2.7.

In the following, we deal with some properties of the Gâteaux derivative of the functional $\Phi$ that are needed in our investigation.
Proposition 2.3. The functional $\Phi^{\prime}: W_{0}^{1, \vec{p}}(\Omega) \rightarrow\left(W_{0}^{1, \vec{p}}(\Omega)\right)^{*}$ defined in (2.14) is monotone and coercive.
Proof. First we prove that $\Phi^{\prime}$ is monotone, i.e.

$$
\begin{equation*}
\left\langle\Phi^{\prime}(u)-\Phi^{\prime}(v), u-v\right\rangle \geq 0 \quad \text { for all } u, v \in W_{0}^{1, \vec{p}}(\Omega) \tag{2.15}
\end{equation*}
$$

To this end, we use the following inequality

$$
\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{r-2} \frac{\partial u}{\partial x_{i}}-\left|\frac{\partial v}{\partial x_{i}}\right|^{r-2} \frac{\partial v}{\partial x_{i}}\right) \frac{\partial(u-v)}{\partial x_{i}} \geq C\left|\frac{\partial(u-v)}{\partial x_{i}}\right|^{r}
$$

for each $r \geq 2$ and for some constant $C>0$, see Simon [33] or Lindqvist [22]. Indeed, using the previous inequality we have

$$
\begin{aligned}
& \left\langle\Phi^{\prime}(u)-\Phi^{\prime}(v), u-v\right\rangle \\
& =\sum_{i=1}^{N} \int_{\Omega}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial u}{\partial x_{i}}-\left|\frac{\partial v}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial v}{\partial x_{i}}\right) \frac{\partial(u-v)}{\partial x_{i}} \mathrm{~d} x \\
& \quad+\sum_{i=1}^{N} \int_{\Omega}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial u}{\partial x_{i}}-\left|\frac{\partial v}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial v}{\partial x_{i}}\right) \frac{\partial(u-v)}{\partial x_{i}} \mathrm{~d} x \\
& \geq \sum_{i=1}^{N} \int_{\Omega}\left(C_{0}\left|\frac{\partial(u-v)}{\partial x_{i}}\right|^{p_{i}}+C_{1}\left|\frac{\partial(u-v)}{\partial x_{i}}\right|^{q_{i}}\right) \mathrm{d} x \\
& \geq 0
\end{aligned}
$$

and (2.15) is achieved.
Now, we prove that $\Phi^{\prime}$ is coercive. We observe that

$$
\begin{equation*}
\left\langle\Phi^{\prime}(u), u\right\rangle=\sum_{i=1}^{N} \int_{\Omega}\left(\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}}+\left|\frac{\partial u}{\partial x_{i}}\right|^{q_{i}}\right) \mathrm{d} x \geq \sum_{i=1}^{N} \|\left.\frac{\partial u}{\partial x_{i}}\right|_{L^{p_{i}}(\Omega)} ^{p_{i}} \tag{2.16}
\end{equation*}
$$

Moreover, let $j \in\{1, \ldots, N\}$ be such that

$$
\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{j}}\right|^{p_{j}} \mathrm{~d} x\right)^{\frac{1}{p_{j}}}:=\max _{1 \leq i \leq N}\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x\right)^{\frac{1}{p_{i}}}
$$

Then,

$$
\begin{aligned}
\|u\|_{W_{0}^{1, \vec{p}}(\Omega)} & =\sum_{i=1}^{N}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)} \leq N\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{j}}\right|^{p_{j}} \mathrm{~d} x\right)^{\frac{1}{p_{j}}} \\
& \leq N\left(\int_{\Omega}\left|\frac{\partial u}{\partial x_{1}}\right|^{p_{1}} \mathrm{~d} x+\ldots+\int_{\Omega}\left|\frac{\partial u}{\partial x_{N}}\right|^{p_{i}} \mathrm{~d} x\right)^{\frac{1}{p_{j}}} \\
& =N\left(\sum_{i=1}^{N}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}^{p_{i}}\right)^{\frac{1}{p_{j}}} .
\end{aligned}
$$

Thus, we get

$$
\begin{equation*}
\sum_{i=1}^{N}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}^{p_{i}} \geq \frac{1}{N^{p_{j}}}\|u\|_{W_{0}^{1, \vec{p}}(\Omega)}^{p_{j}} \tag{2.17}
\end{equation*}
$$

From (2.16) and (2.17), we derive

$$
\left\langle\Phi^{\prime}(u), u\right\rangle \geq \frac{1}{N^{p_{j}}}\|u\|_{W_{0}^{1, \vec{p}}(\Omega)}^{p_{j}}
$$

namely

$$
\frac{\left\langle\Phi^{\prime}(u), u\right\rangle}{\|u\|_{W_{0}^{1, \vec{p}}(\Omega)}} \geq \frac{1}{N^{p_{j}}}\|u\|_{W_{0}^{1, \vec{p}}(\Omega)}^{p_{j}-1},
$$

and this implies that $\Phi^{\prime}$ is coercive.
Proposition 2.4. The map $\Phi^{\prime}: W_{0}^{1, \vec{p}}(\Omega) \rightarrow\left(W_{0}^{1, \vec{p}}(\Omega)\right)^{*}$ has the $\left(\mathrm{S}_{+}\right)$-property, that is

$$
\text { if } \quad u_{n} \rightharpoonup u \quad \text { in } W_{0}^{1, \vec{p}}(\Omega) \quad \text { and } \quad \limsup _{n \rightarrow \infty}\left\langle\Phi^{\prime}\left(u_{n}\right), u_{n}-u\right\rangle \leq 0 \text {, }
$$

then $u_{n} \rightarrow u$ in $W_{0}^{1, \vec{p}}(\Omega)$.
Proof. Let $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subset W_{0}^{1, \vec{p}}(\Omega)$ be such that $u_{n} \rightharpoonup u$ in $W_{0}^{1, \vec{p}}(\Omega)$ and

$$
\begin{equation*}
\limsup _{n \rightarrow+\infty}\left\langle\Phi^{\prime}\left(u_{n}\right), u_{n}-u\right\rangle \leq 0 . \tag{2.18}
\end{equation*}
$$

First, we observe that

$$
\begin{equation*}
\lim _{n \rightarrow+\infty}\left\langle\Phi^{\prime}(u), u_{n}-u\right\rangle=0, \tag{2.19}
\end{equation*}
$$

since $\Phi^{\prime}(u)$ is a linear operator in $W_{0}^{1, \vec{p}}(\Omega)$ and $u_{n} \rightharpoonup u$ in $W_{0}^{1, \vec{p}}(\Omega)$. Hence, from (2.18) and (2.19) it follows that

$$
\begin{equation*}
\limsup _{n \rightarrow+\infty}\left\langle\Phi^{\prime}\left(u_{n}\right)-\Phi^{\prime}(u), u_{n}-u\right\rangle \leq 0 . \tag{2.20}
\end{equation*}
$$

Now, for all $i=1, \ldots, N$ and for all $w, v \in W_{0}^{1, \vec{p}}(\Omega)$ we set

$$
\begin{aligned}
\mathcal{A}_{i}^{p_{i}}(w)(v) & =\int_{\Omega}\left|\frac{\partial w}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial w}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x, \\
\mathcal{A}_{i}^{q_{i}}(w)(v) & =\int_{\Omega}\left|\frac{\partial w}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial w}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x, \\
\mathcal{B}_{i}^{p_{i}}(w)(v) & =\frac{1}{\left\|\frac{\partial w}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}^{p_{i}-2}} \mathcal{A}_{i}^{p_{i}}(w)(v), \\
\mathcal{B}_{i}^{q_{i}}(w)(v) & =\frac{1}{\left\|\frac{\partial w}{\partial x_{i}}\right\|_{L^{q_{i}}(\Omega)}^{q_{i}-2}} \mathcal{A}_{i}^{q_{i}}(w)(v),
\end{aligned}
$$

and

$$
\mathcal{B}_{i}(w)(v)=\mathcal{B}_{i}^{p_{i}}(w)(v)+\mathcal{B}_{i}^{q_{i}}(w)(v) .
$$

Then, we can write $\Phi^{\prime}$ as follows

$$
\Phi^{\prime}(w)(v)=\sum_{i=1}^{N}\left(\mathcal{A}_{i}^{p_{i}}(w)(v)+\mathcal{A}_{i}^{q_{i}}(w)(v)\right) \quad \text { for every } w, v \in W_{0}^{1, \vec{p}}(\Omega) .
$$

By (2.20), we get

$$
\begin{equation*}
\limsup _{n \rightarrow+\infty}\left\langle\mathcal{B}_{i}\left(u_{n}\right)-\mathcal{B}_{i}(u), u_{n}-u\right\rangle \leq 0 \quad \text { for all } i=1, \ldots, N . \tag{2.21}
\end{equation*}
$$

Moreover, we have

$$
\begin{aligned}
& \left\langle\mathcal{B}_{i}\left(u_{n}\right)-\mathcal{B}_{i}(u), u_{n}-u\right\rangle \\
& =\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}^{2}+\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{q_{i}}(\Omega)}^{2}+\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}^{2}+\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{q_{i}}(\Omega)}^{2}
\end{aligned}
$$

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$$
\begin{aligned}
& -\left\langle\mathcal{B}_{i}\left(u_{n}\right), u\right\rangle-\left\langle\mathcal{B}_{i}(u), u_{n}\right\rangle \\
= & \left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}-\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}\right)^{2}+\left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}-\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}\right)^{2} \\
& +2\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}+2\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}} \\
& -\left\langle\mathcal{B}_{i}\left(u_{n}\right), u\right\rangle-\left\langle\mathcal{B}_{i}(u), u_{n}\right\rangle \\
= & \left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}-\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}\right)^{2}+\left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}-\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}\right)^{2} \\
& +\left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}-\left\langle\mathcal{B}_{i}^{p_{i}}\left(u_{n}\right), u\right\rangle\right) \\
& +\left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}-\left\langle\mathcal{B}_{i}^{p_{i}}(u), u_{n}\right\rangle\right) \\
& +\left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}-\left\langle\mathcal{B}_{i}^{q_{i}}\left(u_{n}\right), u\right\rangle\right) \\
& +\left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{q_{i}}(\Omega)}-\left\langle\mathcal{B}_{i}^{q_{i}}(u), u_{n}\right\rangle\right) .
\end{aligned}
$$

Also, applying Hölder's inequality one has

$$
\begin{aligned}
\left|\left\langle\mathcal{B}_{i}^{p_{i}}\left(u_{n}\right), u\right\rangle\right| & \leq\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}, \\
\left|\left\langle\mathcal{B}_{i}^{p_{i}}(u), u_{n}\right\rangle\right| & \leq\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}, \\
\left|\left\langle\mathcal{B}_{i}^{q_{i}}\left(u_{n}\right), u\right\rangle\right| & \leq\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}, \\
\left|\left\langle\mathcal{B}_{i}^{q_{i}}(u), u_{n}\right\rangle\right| & \leq\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{q_{i}}(\Omega)}\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{q_{i}}(\Omega)} .
\end{aligned}
$$

Hence, we derive that

$$
\begin{aligned}
& \left\langle\mathcal{B}_{i}\left(u_{n}\right)-\mathcal{B}_{i}(u), u_{n}-u\right\rangle \\
& \geq\left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}(\Omega)}}-\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}\right)^{2}+\left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{q_{i}(\Omega)}}-\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{q_{i}}(\Omega)}\right)^{2} \\
& \geq\left(\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}-\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}\right)^{2}
\end{aligned}
$$

which, taking (2.21) into account, implies that

$$
\lim _{n \rightarrow+\infty}\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}=\left\|\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)} \quad \text { for all } i=1, \ldots, N .
$$

Then, from Papageorgiou-Winkert [26, Proposition 4.1.11], since $L^{p_{i}}(\Omega)$ is uniformly convex, one has

$$
\lim _{n \rightarrow+\infty}\left\|\frac{\partial u_{n}}{\partial x_{i}}-\frac{\partial u}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}=0 \quad \text { for all } i=1, \ldots, N .
$$

Thus, it follows that

$$
\lim _{n \rightarrow+\infty}\left\|u_{n}-u\right\|_{W_{0}^{1, \vec{p}}(\Omega)}=0
$$

and our claim is proved.
Finally, we point out the following result in order to get information on the sign of the solutions of $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$. For this purpose, let

$$
f^{+}(x, t)= \begin{cases}f(x, 0) & \text { if } t<0 \\ f(x, t) & \text { if } t \geq 0\end{cases}
$$

for all $(x, t) \in \Omega \times \mathbb{R}$ and consider the following problem

$$
\begin{aligned}
-\Delta_{\vec{p}} u-\Delta_{\vec{q}} u & =\lambda f^{+}(x, u) & & \text { in } \Omega, \\
u & =0 & & \text { on } \partial \Omega
\end{aligned}
$$

Lemma 2.5. Assume that $f(x, 0) \geq 0$ for a.a. $x \in \Omega$. Then, any weak solution of problem $\left(D_{\lambda, f^{+}}^{\vec{p}, \vec{q}}\right)$ is nonnegative and it is also a weak solution of problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$.

Proof. Let $u_{0}$ be a weak solution of problem $\left(D_{\lambda, f+}^{\vec{p}, \vec{q}}\right)$, namely

$$
\begin{align*}
& \sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial u_{0}}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x+\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial u_{0}}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x  \tag{2.22}\\
& =\lambda \int_{\Omega} f^{+}\left(x, u_{0}\right) v \mathrm{~d} x
\end{align*}
$$

for all $v \in W_{0}^{1, \vec{p}}(\Omega)$.
In order to prove that $u_{0}$ is nonnegative, put $A=\left\{x \in \Omega: u_{0}(x)<0\right\}$ and $u_{0}^{-}=\min \left\{u_{0}, 0\right\}$. Clearly, $u_{0}^{-} \in W_{0}^{1, \vec{p}}(\Omega)$ (see, for example, Papageorgiou-Winkert [26, Corollary 4.5.19]). Choosing $v=u_{0}^{-}$in (2.22), one has

$$
\begin{aligned}
0 & \leq \sum_{i=1}^{N} \int_{A}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x \\
& \leq \sum_{i=1}^{N} \int_{A}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial u_{0}}{\partial x_{i}} \frac{\partial u_{0}^{-}}{\partial x_{i}} \mathrm{~d} x+\sum_{i=1}^{N} \int_{A}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial u_{0}}{\partial x_{i}} \frac{\partial u_{0}^{-}}{\partial x_{i}} \mathrm{~d} x \\
& =\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial u_{0}}{\partial x_{i}} \frac{\partial u_{0}^{-}}{\partial x_{i}} \mathrm{~d} x+\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial u_{0}}{\partial x_{i}} \frac{\partial u_{0}^{-}}{\partial x_{i}} \mathrm{~d} x \\
& =\lambda \int_{\Omega} f^{+}\left(x, u_{0}\right) u_{0}^{-} \mathrm{d} x \\
& =\lambda \int_{A} f(x, 0) u_{0}^{-} \mathrm{d} x \leq 0
\end{aligned}
$$

that is,

$$
\sum_{i=1}^{N} \int_{A}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x=0
$$

Hence, it holds that

$$
\int_{A}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x=0 \quad \text { for all } i=1, \ldots, N
$$

which yields

$$
\int_{\Omega}\left|\frac{\partial u_{0}-}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x=\int_{A}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x+\int_{\Omega \backslash A}\left|\frac{\partial u_{0}-}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x=0
$$

for all $i=1, \ldots, N$. Therefore, we obtain

$$
\left\|u_{0}^{-}\right\|_{W_{0}^{1, \vec{p}}(\Omega)}=\sum_{i=1}^{N}\left(\int_{\Omega}\left|\frac{\partial u_{0}^{-}}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x\right)^{\frac{1}{p_{i}}}=0
$$

so $u_{0}^{-}=0$ in $\Omega$, which means $u_{0} \geq 0$ in $\Omega$.
Now, we prove that $u_{0}$ is a weak solution for problem $\left(D_{\lambda}^{\overrightarrow{,}, \vec{q}}\right)$. Indeed, from (2.22) one has

$$
\begin{aligned}
& \sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{p_{i}-2} \frac{\partial u_{0}}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x+\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{0}}{\partial x_{i}}\right|^{q_{i}-2} \frac{\partial u_{0}}{\partial x_{i}} \frac{\partial v}{\partial x_{i}} \mathrm{~d} x \\
& =\lambda \int_{\Omega} f^{+}\left(x, u_{0}\right) v \mathrm{~d} x=\lambda \int_{\Omega} f\left(x, u_{0}\right) v \mathrm{~d} x
\end{aligned}
$$

for all $v \in W_{0}^{1, \vec{p}}(\Omega)$, and the conclusion is achieved.
The proofs of our main results are based on the following two critical point theorem due to Bonanno-D'Aguì [7, see Theorem 2.1 and Remark 2.2]. First we recall the definition of the Cerami condition.
Definition 2.6. Let $X$ be a Banach space and $X^{*}$ be its topological dual space. Given $I_{\lambda} \in C^{1}(X)$, we say that $I_{\lambda}$ satisfies the Cerami-condition (C-condition for short), if every sequence $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq X$ such that
$\left(\mathrm{C}_{1}\right)\left|I_{\lambda}\left(u_{n}\right)\right| \leq c_{1}$ for some $c_{1}>0$ and for all $n \in \mathbb{N}$,
$\left(\mathrm{C}_{2}\right)\left(1+\left\|u_{n}\right\|_{X}\right) I_{\lambda}^{\prime}\left(u_{n}\right) \rightarrow 0$ in $X^{*}$ as $n \rightarrow \infty$,
admits a strongly convergent subsequence in $X$.
Theorem 2.7. Let $X$ be a real Banach space and let $\Phi, \Psi: X \rightarrow \mathbb{R}$ be two functionals of class $C^{1}$ such that $\inf _{X} \Phi(u)=\Phi(0)=\Psi(0)=0$. Assume that there are $r \in \mathbb{R}$ and $\tilde{u} \in X$, with $0<\Phi(\tilde{u})<r$, such that

$$
\begin{equation*}
\frac{\sup _{\left.\left.u \in \Phi^{-1}(]-\infty, r\right]\right)} \Psi(u)}{r}<\frac{\Psi(\tilde{u})}{\Phi(\tilde{u})} \tag{2.23}
\end{equation*}
$$

and, for each

$$
\lambda \in \tilde{\Lambda}=] \frac{\Phi(\tilde{u})}{\Psi(\tilde{u})}, \frac{r}{\sup _{\left.\left.u \in \Phi^{-1}(]-\infty, r\right]\right)} \Psi(u)}[
$$

the functional $I_{\lambda}=\Phi-\lambda \Psi$ satisfies the C-condition and it is unbounded from below. Moreover, $\Phi$ is supposed to be coercive.

Then, for each $\lambda \in \tilde{\Lambda}$, the functional $I_{\lambda}$ admits at least two nontrivial critical points $u_{\lambda, 1}, u_{\lambda, 2} \in X$ such that $I\left(u_{\lambda, 1}\right)<0<I\left(u_{\lambda, 2}\right)$.

## 3. Main Results

In this section we formulate and prove our main results. We suppose the following assumptions on the nonlinearity.
$\left(\mathrm{H}_{\mathrm{f}}\right) f: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a Carathéodory function such that the conditions below are satisfied:
$\left(\mathrm{f}_{1}\right)$ there exist $\alpha<\left(p^{-}\right)^{*}$ and constants $a_{1}, a_{2}>0$ such that

$$
|f(x, t)| \leq a_{1}+a_{2}|t|^{\alpha-1}
$$

for a.a. $x \in \Omega$ and for all $t \in \mathbb{R}$;
$\left(\mathrm{f}_{2}\right)$ if $F(x, s)=\int_{0}^{s} f(x, \xi) \mathrm{d} \xi$, then

$$
\lim _{t \rightarrow \pm \infty} \frac{F(x, t)}{|t|^{p^{+}}}=+\infty
$$

uniformly for a. a. $x \in \Omega$;
$\left(\mathrm{f}_{3}\right)$ there exist $\beta, \gamma \in \mathbb{R}$, with

$$
\min \{\beta, \gamma\} \in\left(\left(\alpha-p^{-}\right) \frac{N}{p^{-}}, \alpha\right)
$$

such that

$$
0<m \leq \liminf _{t \rightarrow+\infty} \frac{f(x, t) t-p^{+} F(x, t)}{|t|^{\beta}}
$$

uniformly for a.a. $x \in \Omega$, and

$$
0<m \leq \liminf _{t \rightarrow-\infty} \frac{f(x, t) t-p^{+} F(x, t)}{|t|^{\gamma}}
$$

uniformly for a.a. $x \in \Omega$.
Remark 3.1. We observe that from hypotheses $\left(\mathrm{f}_{1}\right)$ and $\left(\mathrm{f}_{2}\right)$ it follows that

$$
p^{+}<\alpha<\left(p^{-}\right)^{*}
$$

Furthermore, we emphasize that such $\beta$ and $\gamma$ in $\left(\mathrm{f}_{3}\right)$ exists, since

$$
\left(\alpha-p^{-}\right) \frac{N}{p^{-}}=\alpha \frac{N}{p^{-}}-\left(p^{-}\right)^{*} \frac{N-p^{-}}{p^{-}}<\alpha \frac{N}{p^{-}}-\alpha \frac{N-p^{-}}{p^{-}}=\alpha
$$

Finally, we underline that in this context it is possible to choose two different exponents $\beta$ and $\gamma$ for going to $+\infty$ and $-\infty$, respectively.

In the following we give a preliminary result on the energy functional associated to our problem.

Proposition 3.2. Let hypotheses $(\mathrm{H})$ and $\left(\mathrm{H}_{\mathrm{f}}\right)$ be satisfied. Then the functional $I_{\lambda}: W_{0}^{1, \vec{p}}(\Omega) \rightarrow \mathbb{R}$ satisfies the C -condition for each $\lambda>0$.

Proof. Let $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, \vec{p}}(\Omega)$ be a sequence such that $\left(\mathrm{C}_{1}\right)$ and $\left(\mathrm{C}_{2}\right)$ hold, see Definition 2.6. First, we prove that $\left\{u_{n}\right\}_{n \in \mathbb{N}}$ is bounded in $L^{\beta}(\Omega)$.
From $\left(\mathrm{C}_{2}\right)$, we get

$$
\begin{equation*}
\left|\left\langle\Phi^{\prime}\left(u_{n}\right), v\right\rangle-\lambda \int_{\Omega} f\left(x, u_{n}\right) v \mathrm{~d} x\right| \leq \frac{\varepsilon_{n}\|v\|}{1+\left\|u_{n}\right\|}, \tag{3.1}
\end{equation*}
$$

for all $v \in W_{0}^{1, \vec{p}}(\Omega)$ and with $\varepsilon_{n} \rightarrow 0^{+}$. Fix $n \in \mathbb{N}$ and choose $v=u_{n} \in W_{0}^{1, \vec{p}}(\Omega)$. Substituting in (3.1), we derive

$$
\begin{equation*}
-\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x-\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{q_{i}} \mathrm{~d} x+\lambda \int_{\Omega} f\left(x, u_{n}\right) u_{n} \mathrm{~d} x \leq \varepsilon_{n} \tag{3.2}
\end{equation*}
$$

for all $n \in \mathbb{N}$. From $\left(\mathrm{C}_{1}\right)$ we have

$$
\begin{equation*}
\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x+\frac{p^{+}}{q^{+}} \sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{q_{i}} \mathrm{~d} x-\lambda \int_{\Omega} p^{+} F\left(x, u_{n}\right) \mathrm{d} x \leq p^{+} c_{1} . \tag{3.3}
\end{equation*}
$$

Adding (3.2) and (3.3) and taking into account that $q^{+}<p^{+}$, we obtain

$$
\begin{equation*}
\lambda \int_{\Omega}\left(f\left(x, u_{n}\right) u_{n}-p^{+} F\left(x, u_{n}\right)\right) \mathrm{d} x \leq c_{2} \tag{3.4}
\end{equation*}
$$

for some $c_{2}>0$ and for all $n \in \mathbb{N}$.
Without loss of generality, we may assume $\beta \leq \gamma$. Then, assumptions ( $\mathrm{f}_{1}$ ) and $\left(\mathrm{f}_{3}\right)$ imply that there exist $c_{3} \in(0, m)$ and $c_{4}>0$ such that

$$
c_{3}|s|^{\beta}-c_{4} \leq f(x, s) s-p^{+} F(x, s)
$$

for a.a. $x \in \Omega$ and for all $s \in \mathbb{R}$. Exploiting this in (3.4), we derive

$$
\left\|u_{n}\right\|_{L^{\beta}(\Omega)}^{\beta} \leq c_{5} \quad \text { for some } c_{5}>0 \text { and for all } n \in \mathbb{N}
$$

namely $\left\{u_{n}\right\}_{n \in \mathbb{N}}$ is bounded in $L^{\beta}(\Omega)$.
Now, we prove that $\left\{u_{n}\right\}_{n \in \mathbb{N}}$ is bounded in $W_{0}^{1, \vec{p}}(\Omega)$.
From hypotheses $\left(\mathrm{f}_{1}\right)$ and $\left(\mathrm{f}_{3}\right)$ it follows that $\beta<\alpha<\left(p^{-}\right)^{*}$. Hence, there exists $t \in(0,1)$ such that

$$
\begin{equation*}
\frac{1}{\alpha}=\frac{1-t}{\beta}+\frac{t}{\left(p^{-}\right)^{*}} \tag{3.5}
\end{equation*}
$$

By the interpolation inequality (see Papageorgiou-Winkert [26, Proposition 2.3.17]), we get

$$
\left\|u_{n}\right\|_{L^{\alpha}(\Omega)} \leq\left\|u_{n}\right\|_{L^{\beta}(\Omega)}^{1-t}\left\|u_{n}\right\|_{L^{\left(p^{-}\right)^{*}}(\Omega)}^{t} \quad \text { for all } n \in \mathbb{N}
$$

Since $\left\{u_{n}\right\}_{n \in \mathbb{N}}$ is bounded in $L^{\beta}(\Omega)$, using also Proposition 2.2, one has

$$
\begin{equation*}
\left\|u_{n}\right\|_{L^{\alpha}(\Omega)}^{\alpha} \leq c_{6}\left\|u_{n}\right\|_{W_{0}^{1, \vec{p}}(\Omega)}^{t \alpha} \quad \text { for all } n \in \mathbb{N} \tag{3.6}
\end{equation*}
$$

with some $c_{6}>0$. Choosing $v=u_{n} \in W_{0}^{1, \vec{p}}(\Omega)$ in (3.1), we have

$$
\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{p_{i}} \mathrm{~d} x+\sum_{i=1}^{N} \int_{\Omega}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{q_{i}} \mathrm{~d} x-\lambda \int_{\Omega} f\left(x, u_{n}\right) u_{n} \mathrm{~d} x \leq \varepsilon_{n}
$$

for all $n \in \mathbb{N}$. By (2.17), $\left(\mathrm{f}_{1}\right)$ and (3.6) we obtain that there exists $j \in\{1, \ldots, N\}$ such that

$$
\begin{align*}
\frac{1}{N^{p_{j}}}\left\|u_{n}\right\|_{W_{0}^{1, \vec{p}}(\Omega)}^{p_{j}} & \leq \sum_{i=1}^{N}\left\|\frac{\partial u_{n}}{\partial x_{i}}\right\|_{L^{p_{i}}(\Omega)}^{p_{i}} \\
& \leq \lambda \int_{\Omega} f\left(x, u_{n}\right) u_{n} \mathrm{~d} x+\varepsilon_{n}  \tag{3.7}\\
& \leq \lambda c_{7}\left[1+\left\|u_{n}\right\|_{W_{0}^{1, \vec{p}}(\Omega)}^{t \alpha}\right]+\varepsilon_{n}
\end{align*}
$$

for some $c_{7}>0$ and for all $n \in \mathbb{N}$.
From (3.5) and ( $\mathrm{f}_{3}$ ) follows that $t \alpha<p_{j}$, indeed

$$
\begin{align*}
t \alpha=\frac{\left(p^{-}\right)^{*}(\alpha-\beta)}{\left(p^{-}\right)^{*}-\beta} & =\frac{N p^{-}(\alpha-\beta)}{N p^{-}-N \beta+\beta p^{-}} \\
& <\frac{N p^{-}(\alpha-\beta)}{N p^{-}-N \beta+\left(\alpha-p^{-}\right) \frac{N}{p^{-}} p^{-}}=p^{-}  \tag{3.8}\\
& \leq p_{j} .
\end{align*}
$$

Then, (3.7) and (3.8) imply that $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, \vec{p}}(\Omega)$ is bounded.
Finally, we prove that $\left\{u_{n}\right\}_{n \in \mathbb{N}}$ admits a strongly convergent subsequence in $W_{0}^{1, \vec{p}}(\Omega)$. Because of the boundedness of $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, \vec{p}}(\Omega)$, there exists a subsequence, not relabeled, such that

$$
\begin{equation*}
u_{n} \rightharpoonup u \quad \text { in } W_{0}^{1, \vec{p}}(\Omega) \quad \text { and } \quad u_{n} \rightarrow u \quad \text { in } L^{\alpha}(\Omega) \tag{3.9}
\end{equation*}
$$

We choose $v=u_{n}-u \in W_{0}^{1, \vec{p}}(\Omega)$ in (3.1). Passing to the limit as $n \rightarrow \infty$ and using (3.9), we derive

$$
\lim _{n \rightarrow \infty}\left\langle\Phi^{\prime}\left(u_{n}\right), u_{n}-u\right\rangle=0
$$

Since $\Phi^{\prime}$ has the ( $\mathrm{S}_{+}$)-property (see Proposition 2.4), it follows that $u_{n} \rightarrow u$ in $W_{0}^{1, \vec{p}}(\Omega)$ and this completes the proof.

Now, we present our main result. To this aim, put

$$
\begin{equation*}
R:=\sup _{x \in \Omega} \operatorname{dist}(x, \partial \Omega) \tag{3.10}
\end{equation*}
$$

Standard computations show that there exists $x_{0} \in \Omega$ such that $B\left(x_{0}, R\right) \subseteq \Omega$ and we denote by

$$
\omega_{R}:=\left|B\left(x_{0}, R\right)\right|=\frac{\pi^{\frac{N}{2}}}{\Gamma\left(1+\frac{N}{2}\right)} R^{N}
$$

the measure of the $N$-dimensional ball of radius $R$. Finally, we set

$$
\begin{align*}
\mathcal{K} & =\max _{1 \leq i \leq N}\left\{\frac{1}{R^{p_{i}}}, \frac{1}{R^{q_{i}}}\right\} \frac{2^{N}-1}{q^{-}\left(2^{N-p^{+}}\right)} \omega_{R}  \tag{3.11}\\
\delta & =\max \left\{r^{\frac{1}{p^{-}}}, r^{\frac{1}{p^{+}}}\right\} \sum_{i=1}^{N} p_{i}^{\frac{1}{p_{i}}} \tag{3.12}
\end{align*}
$$

Theorem 3.3. Let hypotheses $(\mathrm{H})$ and $\left(\mathrm{H}_{\mathrm{f}}\right)$ be satisfied. Assume that there exist two constants $r, d>0$ such that
$\left(\mathrm{h}_{1}\right) \mathcal{K} \sum_{i=1}^{N}\left(d^{p_{i}}+d^{q_{i}}\right)<r$,
$\left(\mathrm{h}_{2}\right) F(x, s) \geq 0 \quad$ for a.a. $x \in \Omega$ and for all $s \in[0, d]$,
$\left(\mathrm{h}_{3}\right) \frac{a_{1} T_{1} \delta+\frac{a_{2}}{\alpha} T_{\alpha}^{\alpha} \delta^{\alpha}}{r}<\frac{1}{\mathcal{K}} \frac{\int_{B\left(x_{0}, \frac{R}{2}\right)} F(x, d) \mathrm{d} x}{\sum_{i=1}^{N}\left(d^{p_{i}}+d^{q_{i}}\right)}$,
where $a_{1}, a_{2}, \alpha$ are given in $\left(f_{1}\right), T_{1}$ and $T_{\alpha}$ are defined by formula (2.10), $R$ is given in (3.10), and $\mathcal{K}$ as well as $\delta$ are defined by (3.11) and (3.12), respectively.

Then, for each

$$
\lambda \in \Lambda:=] \mathcal{K} \frac{\sum_{i=1}^{N}\left(d^{p_{i}}+d^{q_{i}}\right)}{\int_{B\left(x_{0}, \frac{R}{2}\right)} F(x, d) \mathrm{d} x}, \frac{r}{a_{1} T_{1} \delta+\frac{a_{2}}{\alpha} T_{\alpha}^{\alpha} \delta^{\alpha}}[,
$$

problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ has at least two nontrivial weak solutions $u_{\lambda}, v_{\lambda} \in W_{0}^{\vec{p}}(\Omega)$ such that $I_{\lambda}\left(u_{\lambda}\right)<0<I_{\lambda}\left(v_{\lambda}\right)$.
Proof. Our aim is to apply Theorem 2.7 for $X=W_{0}^{1, \vec{p}}(\Omega)$ and $\Phi, \Psi$ defined as in (2.13). The functionals $\Phi$ and $\Psi$ satisfy all the required regularity properties, since $\Phi$ is coercive by construction (see Proposition 2.3), the energy functional $I_{\lambda}$ satisfies the C-condition due to Proposition 3.2 and it is unbounded from below by ( $\mathrm{f}_{2}$ ) and finally

$$
\inf _{u \in W_{0}^{1, \vec{p}}(\Omega)} \Psi(u)=\Psi(0)=\Phi(0)
$$

Moreover, the interval $\Lambda$ is nonempty because of assumption $\left(h_{3}\right)$. Thus, it remains to verify hypothesis (2.23). First, we observe that

$$
\begin{equation*}
\Phi^{-1}(]-\infty, r[) \subseteq\left\{u \in W_{0}^{1, \vec{p}}(\Omega):\|u\|_{W_{0}^{1, \vec{p}}(\Omega)} \leq \delta\right\} \tag{3.13}
\end{equation*}
$$

Then, from $\left(f_{1}\right),(2.9)$ and (3.13) we estimate that

$$
\begin{align*}
& \frac{\sup _{\left.\left.u \in \Phi^{-1}(]-\infty, r\right]\right)} \Psi(u)}{r} \\
& \leq \frac{\sup _{\left.\left.u \in \Phi^{-1}(]-\infty, r\right]\right)}\left(a_{1}\|u\|_{L^{1}(\Omega)}+\frac{a_{2}}{\alpha}\|u\|_{L^{\alpha}(\Omega)}^{\alpha}\right)}{r}  \tag{3.14}\\
& \leq \frac{\sup _{\left.\left.u \in \Phi^{-1}(]-\infty, r\right]\right)}\left[a_{1} T_{1}\|u\|_{W_{0}^{1, \vec{p}}(\Omega)}+\frac{a_{2}}{\alpha} T_{\alpha}^{\alpha}\|u\|_{W_{0}^{1, \vec{p}}(\Omega)}^{\alpha}\right]}{r} \\
& \leq \frac{a_{1} T_{1} \delta+\frac{a_{2}}{\alpha} T_{\alpha}^{\alpha} \delta^{\alpha}}{r} .
\end{align*}
$$

On the other hand, we introduce the following function

$$
\tilde{u}(x)= \begin{cases}0 & \text { if } x \in \Omega \backslash B\left(x_{0}, R\right) \\ \frac{2 d}{R}\left(R-\left|x-x_{0}\right|\right) & \text { if } B\left(x_{0}, R\right) \backslash B\left(x_{0}, \frac{R}{2}\right), \\ d & \text { if } x \in B\left(x_{0}, \frac{R}{2}\right),\end{cases}
$$

where $R$ is given in (3.10). Clearly, $\tilde{u} \in W_{0}^{1, \vec{p}}(\Omega)$. From assumption $\left(\mathrm{h}_{2}\right)$, we get

$$
\begin{align*}
\Psi(\tilde{u})= & \int_{B\left(x_{0}, R\right) \backslash B\left(x_{0}, \frac{R}{2}\right)} F\left(x, \frac{2 d}{R}\left(R-\left|x-x_{0}\right|\right)\right) \mathrm{d} x \\
& +\int_{B\left(x_{0}, \frac{R}{2}\right)} F(x, d) \mathrm{d} x  \tag{3.15}\\
\geq & \int_{B\left(x_{0}, \frac{R}{2}\right)} F(x, d) \mathrm{d} x .
\end{align*}
$$

Furthermore, it holds

$$
\begin{align*}
\Phi(\tilde{u})= & \sum_{i=1}^{N} \frac{1}{p_{i}} \int_{B\left(x_{0}, R\right) \backslash B\left(x_{0}, \frac{R}{2}\right)}\left(\frac{2 d}{R}\right)^{p_{i}} \mathrm{~d} x \\
& +\sum_{i=1}^{N} \frac{1}{q_{i}} \int_{B\left(x_{0}, R\right) \backslash B\left(x_{0}, \frac{R}{2}\right)}\left(\frac{2 d}{R}\right)^{q_{i}} \mathrm{~d} x  \tag{3.16}\\
\leq & \frac{2^{p^{+}}}{q^{-}} \sum_{i=1}^{N}\left(d^{p_{i}}+d^{q_{i}}\right) \max _{1 \leq i \leq N}\left\{\frac{1}{R^{p_{i}}}, \frac{1}{R^{q_{i}}}\right\}\left[\frac{2^{N}-1}{2^{N}} \omega_{R}\right] \\
= & \mathcal{K} \sum_{i=1}^{N}\left(d^{p_{i}}+d^{q_{i}}\right) .
\end{align*}
$$

Hypothesis $\left(h_{1}\right)$ implies that $0<\Phi(\tilde{u})<r$ and combining (3.14), (3.15) as well as (3.16) we have

$$
\begin{aligned}
\sup _{\left.\left.u \in \Phi^{-1}(]-\infty, r\right]\right)} \Psi(u) & \leq \frac{a_{1} T_{1} \delta+\frac{a_{2}}{\alpha} T_{\alpha}^{\alpha} \delta^{\alpha}}{r} \\
& <\frac{1}{\mathcal{K}} \frac{\int_{B\left(x_{0}, \frac{R}{2}\right)} F(x, d) \mathrm{d} x}{\sum_{i=1}^{N}\left(d^{p_{i}}+d^{q_{i}}\right)} \\
& \leq \frac{\Psi(\tilde{u})}{\Phi(\tilde{u})}
\end{aligned}
$$

This proves (2.23). Hence, since $\Lambda \subseteq \tilde{\Lambda}$, Theorem 2.7 ensures that problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ admits at least two nontrivial weak solutions $u_{\lambda}, v_{\lambda} \in W_{0}^{1, \vec{p}}(\Omega)$ with opposite energy sign.

The following result is a consequence of Lemma 2.5 and of Theorem 3.3.
Theorem 3.4. Let hypotheses $(\mathrm{H})$ and $\left(\mathrm{H}_{\mathrm{f}}\right)$ be satisfied. Assume $\left(\mathrm{h}_{1}\right)-\left(\mathrm{h}_{3}\right)$ and that $f(x, 0) \geq 0$ for a.a. $x \in \Omega$. Then, for every

$$
\lambda \in \Lambda:=] \mathcal{K} \frac{\sum_{i=1}^{N}\left(d^{p_{i}}+d^{q_{i}}\right)}{\int_{B\left(x_{0}, \frac{R}{2}\right)} F(x, d) \mathrm{d} x}, \frac{r}{a_{1} T_{1} \delta+\frac{a_{2}}{\alpha} T_{\alpha}^{\alpha} \delta^{\alpha}}[
$$

problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ has at least two nonnegative weak solutions $u_{\lambda}, v_{\lambda} \in W_{0}^{\vec{p}}(\Omega)$ such that $I_{\lambda}\left(u_{\lambda}\right)<0<I_{\lambda}\left(v_{\lambda}\right)$.

Now, we consider the particular case in which the nonlinearity is nonnegative.

Theorem 3.5. Let hypotheses $(\mathrm{H})$ and $\left(\mathrm{H}_{\mathrm{f}}\right)$ be satisfied. Assume that $f$ is nonnegative and

$$
\left(\mathrm{h}_{4}\right) \limsup _{t \rightarrow 0^{+}} \frac{\inf _{x \in \Omega} F(x, t)}{t^{q^{-}}}=+\infty
$$

Then, for each $\lambda \in] 0, \lambda^{*}[$, with

$$
\begin{equation*}
\lambda^{*}=\sup _{r>0} \frac{r}{a_{1} T_{1} \delta+\frac{a_{2}}{\alpha} T_{\alpha}^{\alpha} \delta^{\alpha}} \tag{3.17}
\end{equation*}
$$

problem $\left(D_{\lambda}^{\vec{p}, \vec{q}}\right)$ admits at least two nonnegative weak solutions $u_{\lambda}, v_{\lambda} \in W_{0}^{1, \vec{p}}(\Omega)$ such that $I_{\lambda}\left(u_{\lambda}\right)<0<I_{\lambda}\left(v_{\lambda}\right)$.

Proof. Fix $\lambda \in] 0, \lambda^{*}[$, then there exists $r>0$ such that

$$
\lambda<\frac{r}{a_{1} T_{1} \delta+\frac{a_{2}}{\alpha} T_{\alpha}^{\alpha} \delta^{\alpha}} .
$$

From $\left(h_{4}\right)$ follows that

$$
\limsup _{t \rightarrow 0^{+}} \frac{1}{\mathcal{K}} \frac{\int_{B\left(x_{0}, \frac{R}{2}\right)} F(x, t) \mathrm{d} x}{\sum_{i=1}^{N}\left(t^{p_{i}}+t^{q_{i}}\right)} \geq \frac{\omega_{\frac{R}{2}}}{\mathcal{K}} \limsup _{t \rightarrow 0^{+}} \frac{\inf _{x \in \Omega} F(x, t)}{\sum_{i=1}^{N}\left(t^{p_{i}}+t^{q_{i}}\right)}=+\infty
$$

since

$$
\frac{1}{\sum_{i=1}^{N}\left(t^{p_{i}}+t^{q_{i}}\right)} \geq \frac{1}{2 N t^{q^{-}}}
$$

Then, in correspondence of $\frac{1}{\lambda}$, there exists $t>0$ small enough such that

$$
\frac{1}{\mathcal{K}} \frac{\int_{B\left(x_{0}, \frac{R}{2}\right)} F(x, t) \mathrm{d} x}{\sum_{i=1}^{N}\left(t^{p_{i}}+t^{q_{i}}\right)}>\frac{1}{\lambda}>\frac{a_{1} T_{1} \delta+\frac{a_{2}}{\alpha} T_{\alpha}^{\alpha} \delta^{\alpha}}{r}
$$

namely assumption $\left(h_{3}\right)$ is satisfied. Since $\left(h_{2}\right)$ follows from the sign assumption on the nonlinearity, we can apply Theorem 3.3 and Lemma 2.5 to complete the proof.

Finally, we deal with the autonomous case and we present an existence result which is a consequence of Theorem 3.5. Consider the autonomous problem

$$
\begin{aligned}
-\Delta_{\vec{p}} u-\Delta_{\vec{q}} u & =\lambda g(u) & & \text { in } \Omega, \\
u & =0 & & \text { on } \partial \Omega,
\end{aligned}
$$

$$
\left(A D_{\lambda}^{\vec{p}, \vec{q}}\right)
$$

where $g: \mathbb{R} \rightarrow \mathbb{R}$ is a nonnegative continuous function. From Lemma 2.5 it follows that we can consider the nonlinearity only in $[0,+\infty)$. We assume the following:
$\left(\mathrm{H}_{\mathrm{g}}\right)\left(\mathrm{g}_{1}\right)$ there exist $\alpha<\left(p^{-}\right)^{*}$ and constants $a_{1}, a_{2}>0$ such that

$$
g(t) \leq a_{1}+a_{2}|t|^{\alpha-1}
$$

for all $t \geq 0$;
$\left(\mathrm{g}_{2}\right)$ if $G(s)=\int_{0}^{s} g(\xi) \mathrm{d} \xi$, then

$$
\lim _{t \rightarrow+\infty} \frac{G(t)}{t^{p^{+}}}=+\infty
$$

( $\mathrm{g}_{3}$ ) there exists $\beta \in \mathbb{R}$, with

$$
\beta \in\left(\left(\alpha-p^{-}\right) \frac{N}{p^{-}}, \alpha\right)
$$

such that

$$
0<m \leq \liminf _{t \rightarrow+\infty} \frac{g(t) t-p^{+} G(t)}{t^{\beta}}
$$

The following result holds.
Corollary 3.6. Let hypotheses $(\mathrm{H})$ and $\left(\mathrm{H}_{\mathrm{g}}\right)$ be satisfied. Assume that

$$
\left(\mathrm{h}_{4}{ }^{\prime}\right)
$$

$$
\limsup _{t \rightarrow 0^{+}} \frac{G(t)}{t^{q^{-}}}=+\infty
$$

Then, for each $\lambda \in] 0, \lambda^{*}\left[\right.$, with $\lambda^{*}$ defined in (3.17), the problem $\left(A D_{\lambda}^{\vec{p}, \vec{q}}\right)$ admits at least two nonnegative weak solutions $u_{\lambda}, v_{\lambda} \in W_{0}^{1, \vec{p}}(\Omega)$ such that $I_{\lambda}\left(u_{\lambda}\right)<0<$ $I_{\lambda}\left(v_{\lambda}\right)$.

In conclusion, we provide an example.
Example 3.7. Consider two constants $c, \kappa$ such that

$$
c \geq 1, \quad p^{+}<\kappa<\left(p^{-}\right)^{*} \quad \text { and } \quad \frac{\kappa}{p^{-}}-\frac{\kappa}{N}<1
$$

Let $g:[0,+\infty) \rightarrow \mathbb{R}$ be a function defined by

$$
g(t)=(t+c)^{\kappa-1}(\kappa \log (t+c)+1) \quad \text { for all } t \geq 0
$$

Then, $g$ satisfies assumptions $\left(\mathrm{H}_{\mathrm{g}}\right)$ with $\beta=\kappa$ and $\alpha=\kappa+\sigma$, with $\sigma>0$ small enough such that

$$
\alpha<\left(p^{-}\right)^{*} \quad \text { and } \quad \frac{\alpha}{p^{-}}-\frac{\kappa}{N}<1
$$

Moreover, the function $g$ satisfies assumption ( $\mathrm{h}_{4}{ }^{\prime}$ ), hence we can apply Corollary 3.6 to get the existence of two nonnegative weak solutions of problem $\left(A D_{\lambda}^{\vec{p}, \vec{q}}\right)$ with opposite energy sign.

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