# SEQUENCES OF NODAL SOLUTIONS FOR CRITICAL DOUBLE PHASE PROBLEMS WITH VARIABLE EXPONENTS 

NIKOLAOS S. PAPAGEORGIOU, FRANCESCA VETRO, AND PATRICK WINKERT


#### Abstract

In this paper, we study a double phase problem with both variable exponents. Such problem has a reaction consisting of a Carathéodory perturbation defined only locally and of a critical term. The presence of the critical term does not permit to use results of the critical point theory for the corresponding energy functional. Consequently, using suitable cut-off functions and truncation techniques we focus on an auxiliary coercive problem on which, differently from our main problem, we can act with variational tools. In this way, we are able to produce a sequence of sign-changing solutions to our main problem converging to 0 in $L^{\infty}$ and in the Musielak-Orlicz Sobolev space.


## 1. Introduction

Let $\Omega \subseteq \mathbb{R}^{N}$ with $N \geq 2$ be a bounded domain with Lipschitz boundary $\partial \Omega$. Given $r \in C(\Omega)$, we define

$$
r^{-}=\min _{x \in \bar{\Omega}} r(x) \quad \text { and } \quad r^{+}=\max _{x \in \bar{\Omega}} r(x) .
$$

In this paper, we focus on the following critical double phase Dirichlet problem

$$
\begin{align*}
-\operatorname{div}\left(|\nabla u|^{p(x)-2} \nabla u+\mu(x)|\nabla u|^{q(x)-2} \nabla u\right) & =f(x, u)+|u|^{p^{*}(x)-2} u & & \text { in } \Omega  \tag{1.1}\\
u & =0 & & \text { on } \partial \Omega .
\end{align*}
$$

Here, we suppose the following hypotheses on the exponents, the weight function and the perturbation term:
(H1) $p, q \in C(\bar{\Omega})$ are such that $1<p(x)<N, p(x)<q(x)<p^{*}(x):=\frac{N p(x)}{N-p(x)}$ for all $x \in \bar{\Omega}$ and $0 \leq \mu(\cdot) \in L^{\infty}(\Omega) \backslash\{0\}$.
(H2) $f: \Omega \times\left[-\eta_{0}, \eta_{0}\right] \rightarrow \mathbb{R}$, with $\eta_{0}>0$, is a Carathéodory function such that $f(x, \cdot)$ is odd for a.a. $x \in \Omega$ and
(i) there exists $a_{0} \in L^{\infty}(\Omega)$ such that

$$
|f(x, s)| \leq a_{0}(x) \quad \text { for a.a. } x \in \Omega \text { and for all }|s| \leq \eta_{0}
$$

(ii) there exist $\delta \in\left(0, \min \left\{\frac{\eta_{0}}{2}, 1\right\}\right)$ and $\tau \in C(\bar{\Omega})$ with $1 \leq \tau(x) \leq \tau^{+}<p^{-}$ such that

$$
c_{0}|s|^{\tau(x)} \leq f(x, s) s \leq \tilde{c}_{0}|s|^{\tau(x)}
$$

for some $c_{0}>\frac{2 \tau^{+}}{p^{-}}$and $\tilde{c}_{0}>c_{0}$, for a.a. $x \in \Omega$ and for all $|s| \leq \delta$.

[^0]So, in the right-hand side of problem (1.1) we find the combined effects of a Carathéodory perturbation $f(x, \cdot)$ which is defined only locally and of a critical term $u \rightarrow|u|^{p^{*}(x)-2} u$ with $p^{*}(\cdot)$ being the critical exponent corresponding to $p(\cdot)$.

A function $u \in W_{0}^{1, \mathcal{H}}(\Omega)$ (the Musielak-Orlicz Sobolev space, see Section 2) is said to be a weak solution of problem (1.1) if

$$
\int_{\Omega}\left(|\nabla u|^{p(x)-2} \nabla u+\mu(x)|\nabla u|^{q(x)-2} \nabla u\right) \cdot \nabla h \mathrm{~d} x=\int_{\Omega}\left(f(x, u)+|u|^{p^{*}(x)-2} u\right) h \mathrm{~d} x
$$

is satisfied for all $h \in W_{0}^{1, \mathcal{H}}(\Omega)$.
We point out that the main goal of this paper is to produce nodal (that is, signchanging) solutions for problem (1.1). Precisely, we here establish the following result.

Theorem 1.1. Let hypotheses (H1) and (H2) be satisfied. Then, problem (1.1) has a sequence

$$
\left\{z_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, \mathcal{H}}(\Omega) \cap L^{\infty}(\Omega)
$$

of nodal (that is, sign-changing) solutions such that

$$
\left\|z_{n}\right\| \rightarrow 0 \quad \text { and } \quad\left\|z_{n}\right\|_{\infty} \rightarrow 0 \quad \text { as } n \rightarrow+\infty
$$

The proof of Theorem 1.1 is based on the use of cut-off functions, truncation techniques and also a generalized version of the symmetric mountain pass theorem due to Kajikiya [19, Theorem 1]. Hence, using suitable cut-off functions and truncation techniques we introduce an auxiliary coercive problem (see problem (3.4), Section 3) on which differently from problem (1.1) we can act with variational tools. In fact, we point out that the presence of the critical term $u \rightarrow|u|^{p^{*}(x)-2} u$ in the right-hand side of problem (1.1) does not permit us to apply results of the critical point theory to the corresponding energy functional. Therefore, we focus on the auxiliary coercive problem (3.4). We show the existence of extremal constant sign solutions for such problem (see again Section 3). After that, using these extremal solutions and Kajikiya's theorem we are able to produce a sequence of sign-changing solutions for problem (1.1) (see Section 4).

Also, we remark that our work was inspired by a recent paper of Liu-Papageorgiou [23]. Similar to our finding here, the authors in [23] consider a double phase Dirichlet problem exhibiting the combined effects of a Carathéodory perturbation defined only locally and of a critical term, but they work with constant exponents and under more restrictive conditions. Thus, we extend the results of Liu-Papageorgiou [23] to the case of a double phase operator with both variable exponents and under weaker conditions. Indeed, we are able to skip condition H1 (iii) in [23]. In addition, we point out that our main result extends the one of Papageorgiou-Vetro-Winkert [31] to the case of two variable exponents. In [31] $p \equiv p(x)$ has to be a constant.

Finally, we mention that functionals of type

$$
\mathcal{F}(\omega)=\int_{\Omega}\left(|\nabla \omega|^{p}+\mu(x)|\nabla \omega|^{q}\right) \mathrm{d} x, \quad 1<p<q<N
$$

were studied by Marcellini [26] and Zhikov [40] in order to describe strongly anisotropic materials in the context of homogenization and elasticity. Such functionals also
find application in the study of duality theory and of the Lavrentiev gap phenomenon, see Zhikov [41, 42], and in the context of problems of the calculus of variations, see Marcellini [25, 26].

A first mathematical framework for such functionals has been provided by Baroni-Colombo-Mingione [4], see also the related works by the same authors in [5, 6] and of De Filippis-Mingione [10] about nonautonomous integrals.

So far, there are only few results involving the variable exponent double phase operator. We refer to the recent results of Aberqi-Bennouna-Benslimane-Ragusa [1] for existence results in complete manifolds, Albalawi-Alharthi-Vetro [2] for convection problems with $(p(\cdot), q(\cdot))$-Laplace type operator, Bahrouni-Rădulescu-Winkert [3] for problems with Baouendi-Grushin type operator, Crespo-Blanco-Gasiński-Harjulehto-Winkert [8] for double phase convection problems, Kim-Kim-Oh-Zeng [20] for concave-convex-type double phase problems, Leonardi-Papageorgiou [21] for concave-convex problems, Liu-Pucci [24] for problems without supposing the Ambrosetti-Rabinowitz condition, Vetro-Winkert [35] for parametric problems involving superlinear nonlinearities and Zeng-Rădulescu-Winkert [39] for multivalued problems, see also the references therein. We also recall the papers of ColasuonnoSquassina [7] for eigenvalue problems of double phase type, Farkas-Winkert [12] for Finsler double phase problems, Gasiński-Papageorgiou [13] for locally Lipschitz right-hand sides, Gasiński-Winkert [14, 15] for convection problems and constant sign-solutions, Liu-Dai [22] for a Nehari manifold approach, Papageorgiou-Vetro [29] for superlinear problems, Papageorgiou-Vetro-Vetro [30] for parametric Robin problems, Perera-Squassina [33] for Morse theoretical approach, Vetro-Winkert [36] for parametric convective problems, Vetro-Winkert [37] for critical Robin double phase problems with one variable exponent and Zeng-Bai-Gasiński-Winkert [38] for implicit obstacle problems with multivalued operators.

## 2. Mathematical background

Given a bounded domain $\Omega \subseteq \mathbb{R}^{N}(N \geq 2)$ with Lipschitz boundary $\partial \Omega$, we denote by $M(\Omega)$ the set of all measurable functions $u: \Omega \rightarrow \mathbb{R}$. Let $r \in C(\bar{\Omega})$ be such that $r(x)>1$ for all $x \in \bar{\Omega}$. Then, the usual variable exponent Lebesgue space $L^{r(\cdot)}(\Omega)$ is defined by

$$
L^{r(\cdot)}(\Omega)=\left\{u \in M(\Omega): \rho_{r(\cdot)}(u):=\int_{\Omega}|u|^{r(x)} \mathrm{d} x<+\infty\right\}
$$

We furnish this space with the Luxemburg norm

$$
\|u\|_{r(\cdot)}=\inf \left\{\alpha>0: \rho_{r(\cdot)}\left(\frac{u}{\alpha}\right) \leq 1\right\}
$$

We write $W^{1, r(\cdot)}(\Omega)$ and $W_{0}^{1, r(\cdot)}(\Omega)$ to denote the corresponding Sobolev spaces equipped with the norms $\|\cdot\|_{1, r(\cdot)}$ and $\|\nabla \cdot\|_{r(\cdot)}$, respectively, where $\|\cdot\|_{1, r(\cdot)}$ is given by

$$
\|u\|_{1, r(\cdot)}=\|u\|_{r(\cdot)}+\|\nabla u\|_{r(\cdot)}
$$

see also Diening-Harjulehto-Hästö-Růžička [11] and Harjulehto-Hästö [17]. Below, we recall the relations between $\|\cdot\|_{r(\cdot)}$ and $\rho_{r(\cdot)}$, see again [11].

Proposition 2.1. Let $r \in C(\bar{\Omega})$ be such that $r(x)>1$ for all $x \in \bar{\Omega}$. Then the following hold:
(i) $\|u\|_{r(\cdot)}<1$ (resp. $>1,=1$ ) if and only if $\rho_{r(\cdot)}(u)<1$ (resp. $>1,=1$ );
(ii) if $\|u\|_{r(\cdot)}<1$ then $\|u\|_{r(\cdot)}^{r^{+}} \leq \rho_{r(\cdot)}(u) \leq\|u\|_{r(\cdot)}^{r^{-}}$;
(iii) if $\|u\|_{r(\cdot)}>1$ then $\|u\|_{r(\cdot)}^{r^{-}} \leq \rho_{r(\cdot)}(u) \leq\|u\|_{r(\cdot)}^{r^{+}}$;
(iv) $\|u\|_{r(\cdot)} \rightarrow 0$ if and only if $\rho_{r(\cdot)}(u) \rightarrow 0$;
(v) $\|u\|_{r(\cdot)} \rightarrow+\infty$ if and only if $\rho_{r(\cdot)}(u) \rightarrow+\infty$.

Assume that hypothesis (H1) is satisfied, then we can consider the nonlinear function $\mathcal{H}: \Omega \times[0,+\infty) \rightarrow[0,+\infty)$ defined by

$$
\mathcal{H}(x, t)=t^{p(x)}+\mu(x) t^{q(x)} \quad \text { for all } x \in \Omega \text { and for all } t \geq 0 .
$$

Here, we denote by $\rho_{\mathcal{H}}(\cdot)$ the corresponding modular function, that is,

$$
\rho_{\mathcal{H}}(u)=\int_{\Omega} \mathcal{H}(x,|u|) \mathrm{d} x=\int_{\Omega}\left(|u|^{p(x)}+\mu(x)|u|^{q(x)}\right) \mathrm{d} x .
$$

Using $\rho_{\mathcal{H}}(\cdot)$ we introduce the Musielak-Orlicz space $L^{\mathcal{H}}(\Omega)$ by

$$
L^{\mathcal{H}}(\Omega)=\left\{u \in M(\Omega): \rho_{\mathcal{H}}(u)<+\infty\right\} .
$$

As usual, we equip this space with the Luxemburg norm

$$
\|u\|_{\mathcal{H}}=\inf \left\{\alpha>0: \rho_{\mathcal{H}}\left(\frac{u}{\alpha}\right) \leq 1\right\} .
$$

We underline that also the modular $\rho_{\mathcal{H}}$ and the norm $\|\cdot\|_{\mathcal{H}}$ are related by a close relation, see Crespo-Blanco-Gasiński-Harjulehto-Winkert [8, Proposition 2.13].

Proposition 2.2. Let hypotheses (H1) be satisfied. Then the following hold:
(i) $\|u\|_{\mathcal{H}}<1$ (resp. $>1,=1$ ) if and only if $\rho_{\mathcal{H}}(u)<1$ (resp. $>1,=1$ );
(ii) if $\|u\|_{\mathcal{H}}<1$ then $\|u\|_{\mathcal{H}}^{q^{+}} \leq \rho_{\mathcal{H}}(u) \leq\|u\|_{\mathcal{H}}^{p^{-}}$;
(iii) if $\|u\|_{\mathcal{H}}>1$ then $\|u\|_{\mathcal{H}}^{p^{-}} \leq \rho_{\mathcal{H}}(u) \leq\|u\|_{\mathcal{H}}^{q^{+}}$;
(iv) $\|u\|_{\mathcal{H}} \rightarrow 0$ if and only if $\rho_{\mathcal{H}}(u) \rightarrow 0$;
(v) $\|u\|_{\mathcal{H}} \rightarrow+\infty$ if and only if $\rho_{\mathcal{H}}(u) \rightarrow+\infty$.

Now, starting from the Musielak-Orlicz space $L^{\mathcal{H}}(\Omega)$, we can define the corresponding Musielak-Orlicz Sobolev space $W^{1, \mathcal{H}}(\Omega)$ by

$$
W^{1, \mathcal{H}}(\Omega)=\left\{u \in L^{\mathcal{H}}(\Omega):|\nabla u| \in L^{\mathcal{H}}(\Omega)\right\} .
$$

We consider on this space the norm

$$
\|u\|_{1, \mathcal{H}}=\|u\|_{\mathcal{H}}+\|\nabla u\|_{\mathcal{H}},
$$

where $\|\nabla u\|_{\mathcal{H}}:=\||\nabla u|\|_{\mathcal{H}}$. Also, by $W_{0}^{1, \mathcal{H}}(\Omega)$ we mean the completion of $C_{0}^{\infty}(\Omega)$ in $W^{1, \mathcal{H}}(\Omega)$.

From Crespo-Blanco-Gasiński-Harjulehto-Winkert [8, Proposition 2.12] we know that the spaces $L^{\mathcal{H}}(\Omega), W^{1, \mathcal{H}}(\Omega)$ and $W_{0}^{1, \mathcal{H}}(\Omega)$ are reflexive Banach spaces. Further, on account of Proposition 2.18 in [8] we can endow the space $W_{0}^{1, \mathcal{H}}(\Omega)$ with the equivalent norm

$$
\|u\|=\|\nabla u\|_{\mathcal{H}} \quad \text { for all } u \in W_{0}^{1, \mathcal{H}}(\Omega) .
$$

We recall that the classical Sobolev embedding theorem extends to the space $W_{0}^{1, \mathcal{H}}(\Omega)$ as follows, see Crespo-Blanco-Gasiński-Harjulehto-Winkert [8, Proposition 2.16].
Proposition 2.3. Let hypotheses (H1) be satisfied. Then the following hold:
(i) $W_{0}^{1, \mathcal{H}}(\Omega) \hookrightarrow W_{0}^{1, r(\cdot)}(\Omega)$ is continuous for all $r \in C(\bar{\Omega})$ with $1 \leq r(x) \leq p(x)$ for all $x \in \bar{\Omega}$;
(ii) $W_{0}^{1, \mathcal{H}}(\Omega) \hookrightarrow L^{r(\cdot)}(\Omega)$ is compact for all $r \in C(\bar{\Omega})$ with $1 \leq r(x)<p^{*}(x)$ for all $x \in \bar{\Omega}$.

Finally, for any $s \in \mathbb{R}$ we put $s_{ \pm}=\max \{ \pm s, 0\}$ and we have $s=s_{+}-s_{-}$and $|s|=s_{+}+s_{-}$. In addition, for any function $u: \Omega \rightarrow \mathbb{R}$ we define $u_{ \pm}(\cdot)=[u(\cdot)]_{ \pm}$.

Given a Banach space $X$ and its dual space $X^{*}$, we recall that a functional $\varphi \in C^{1}(X)$ satisfies the Palais-Smale condition (PS-condition for short), if every sequence $\left\{x_{n}\right\}_{n \in \mathbb{N}} \subseteq X$ such that $\left\{\varphi\left(u_{n}\right)\right\}_{n \in \mathbb{N}} \subseteq \mathbb{R}$ is bounded and

$$
\varphi^{\prime}\left(x_{n}\right) \rightarrow 0 \quad \text { in } X^{*} \text { as } n \rightarrow+\infty
$$

admits a strongly convergent subsequence. We denote by $K_{\varphi}$ the set of all critical points of $\varphi$, that is,

$$
K_{\varphi}=\left\{u \in X: \varphi^{\prime}(u)=0\right\}
$$

We also recall that a set $\mathcal{S} \subseteq X$ is said to be downward directed if for given $u_{1}, u_{2} \in \mathcal{S}$ we can find $u \in \mathcal{S}$ such that $u \leq u_{1}$ and $u \leq u_{2}$. Analogously, $\mathcal{S} \subseteq X$ is said to be upward directed if for given $v_{1}, v_{2} \in \mathcal{S}$ we can find $v \in \mathcal{S}$ such that $v_{1} \leq v$ and $v_{2} \leq v$.

## 3. AUXILIARY PROBLEM

In this section, using suitable cut-off functions and truncation techniques, we introduce an auxiliary coercive problem on which we can act with variational tools. The study of this problem helps us to prove the existence of nodal (that is, signchanging) solutions for our main problem (1.1).

Let $\theta \in C^{1}(\mathbb{R})$ be an even cut-off function satisfying the following conditions:

$$
\begin{equation*}
\operatorname{supp} \theta \subseteq\left[-\eta_{0}, \eta_{0}\right], \quad \theta_{\left[\frac{-\eta_{0}}{2}, \frac{\eta_{0}}{2}\right]} \equiv 1 \quad \text { and } \quad 0<\theta \leq 1 \quad \text { on }\left(-\eta_{0}, \eta_{0}\right) \tag{3.1}
\end{equation*}
$$

where $\eta_{0}$ is the positive constant from hypothesis (H2). Using $\theta$, we introduce the Carathéodory function $k: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ defined by

$$
\begin{equation*}
k(x, s)=\theta(s)\left[f(x, s)+|s|^{p^{*}(x)-2} s\right]+(1-\theta(s))|s|^{\tau(x)-2} s \tag{3.2}
\end{equation*}
$$

for all $(x, s) \in \Omega \times \mathbb{R}$, where $\tau$ is given in (H2)(ii). The assumptions on $\theta$ (see (3.1)) and hypothesis (H2) guarantee that

$$
\begin{equation*}
|k(x, s)| \leq c\left(1+|s|^{\tau(x)-1}\right) \tag{3.3}
\end{equation*}
$$

for some $c>0$, for a.a. $x \in \Omega$ and for all $|s| \leq \delta$.
Then, we consider the following auxiliary double phase Dirichlet problem

$$
\begin{align*}
-\operatorname{div}\left(|\nabla u|^{p(x)-2} \nabla u+\mu(x)|\nabla u|^{q(x)-2} \nabla u\right) & =k(x, u) & & \text { in } \Omega  \tag{3.4}\\
u & =0 & & \text { on } \partial \Omega .
\end{align*}
$$

Here, we denote by $\mathcal{S}_{+}$and $\mathcal{S}_{-}$the sets of positive and negative solutions of problem (3.4), respectively. First, we prove that these sets are nonempty.

Proposition 3.1. Let hypotheses (H1) and (H2) be satisfied. Then $\mathcal{S}_{+}$and $\mathcal{S}_{-}$are nonempty subsets in $W_{0}^{1, \mathcal{H}}(\Omega) \cap L^{\infty}(\Omega)$.

Proof. We start by the set $\mathcal{S}_{+}$and show that it is nonempty. To this purpose, we consider the $C^{1}$-functional $\phi_{+}: W_{0}^{1, \mathcal{H}}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\phi_{+}(u)=\int_{\Omega}\left[\frac{1}{p(x)}|\nabla u|^{p(x)}+\frac{\mu(x)}{q(x)}|\nabla u|^{q(x)}\right] \mathrm{d} x-\int_{\Omega} K\left(x, u_{+}\right) \mathrm{d} x
$$

for all $u \in W_{0}^{1, \mathcal{H}}(\Omega)$, where $K(x, s)=\int_{0}^{s} k(x, t) \mathrm{d} t$. Due to (3.3) we obtain

$$
\begin{aligned}
\phi_{+}(u) & \geq \frac{1}{p^{+}} \int_{\Omega}|\nabla u|^{p(x)} \mathrm{d} x+\frac{1}{q^{+}} \int_{\Omega} \mu(x)|\nabla u|^{q(x)} \mathrm{d} x-\int_{\Omega} K\left(x, u_{+}\right) \mathrm{d} x \\
& \geq \frac{1}{q^{+}} \rho_{\mathcal{H}}(|\nabla u|)-\int_{\Omega} K\left(x, u_{+}\right) \mathrm{d} x \\
& \geq \frac{1}{q^{+}} \rho_{\mathcal{H}}(|\nabla u|)-\int_{\Omega} \int_{0}^{u_{+}} c\left(1+|s|^{\tau(x)-1}\right) \mathrm{d} s \\
& \geq \frac{1}{q^{+}} \rho_{\mathcal{H}}(|\nabla u|)-c \int_{\Omega} u_{+} \mathrm{d} x-\frac{c}{\tau^{-}} \int_{\Omega} u_{+}^{\tau(x)} \mathrm{d} x \\
& \geq \frac{1}{q^{+}} \rho_{\mathcal{H}}(|\nabla u|)-c \int_{\Omega}|u| \mathrm{d} x-\frac{c}{\tau^{-}} \int_{\Omega}|u|^{\tau(x)} \mathrm{d} x .
\end{aligned}
$$

We recall that if $\|u\|_{\tau(\cdot)}>1$, due to Proposition 2.1(iii), we have that $\rho_{\tau(\cdot)}(u) \leq$ $\|u\|_{\tau(\cdot)}^{\tau^{+}}$. Similarly, if $\|u\|:=\|\nabla u\|_{\mathcal{H}}>1$ using Proposition 2.2 (iii) we know that $\rho_{\mathcal{H}}(|\nabla u|) \geq\|\nabla u\|_{\mathcal{H}}^{p^{-}}$. According of this, for any $u \in W_{0}^{1, \mathcal{H}}(\Omega)$ such that $\|u\|_{\tau(x)}>1$ and $\|u\|>1$ we can further write

$$
\begin{aligned}
\phi_{+}(u) & \geq \frac{1}{q^{+}}\|u\|^{p^{-}}-c\|u\|_{1}-\frac{c}{\tau^{-}}\|u\|_{\tau(\cdot)}^{\tau^{+}} \\
& \geq \frac{1}{q^{+}}\|u\|^{p^{-}}-c_{1}\|u\|-\frac{c_{1}}{\tau^{-}}\|u\|^{\tau^{+}}
\end{aligned}
$$

for some $c_{1}>0$ since the embeddings $W_{0}^{1, \mathcal{H}}(\Omega) \hookrightarrow L^{1}(\Omega)$ and $W_{0}^{1, \mathcal{H}}(\Omega) \hookrightarrow L^{\tau(\cdot)}(\Omega)$ are compact, see Proposition 2.3(ii). From this, taking into account that $\tau^{+}<p^{-}$ (see (H2)(ii)), we conclude that $\phi_{+}$is coercive. Moreover, using the compactness of the embedding $W_{0}^{1, \mathcal{H}}(\Omega) \hookrightarrow L^{r(\cdot)}(\Omega)$ for any $r \in C(\bar{\Omega})$ with $1 \leq r(x)<p^{*}(x)$ for all $x \in \bar{\Omega}$ (see Proposition 2.3(ii)), we infer that the functional $\phi_{+}$is sequentially weakly lower semicontinuous. Thus, there exists $u_{0} \in W_{0}^{1, \mathcal{H}}(\Omega)$ such that

$$
\phi_{+}\left(u_{0}\right)=\inf \left[\phi_{+}(u): u \in W_{0}^{1, \mathcal{H}}(\Omega)\right]
$$

We show that $u_{0}$ is nontrivial. Hence, we consider a function $\tilde{u} \in \operatorname{int} C(\bar{\Omega})_{+}$and we take $t \in(0,1)$ small enough so that $t \tilde{u}(x) \in(0, \delta]$ for all $x \in \bar{\Omega}$. Then, we have

$$
\begin{align*}
& \phi_{+}(t \tilde{u}) \\
& =\int_{\Omega}\left[\frac{1}{p(x)}|\nabla(t \tilde{u})|^{p(x)}+\frac{\mu(x)}{q(x)}|\nabla(t \tilde{u})|^{q(x)}\right] \mathrm{d} x-\int_{\Omega} K(x, t \tilde{u}) \mathrm{d} x  \tag{3.5}\\
& \leq \frac{t^{p^{-}}}{p^{-}}\left[\int_{\Omega}|\nabla \tilde{u}|^{p(x)} \mathrm{d} x+\int_{\Omega} \mu(x)|\nabla \tilde{u}|^{q(x)} \mathrm{d} x\right]-\int_{\Omega} K(x, t \tilde{u}) \mathrm{d} x
\end{align*}
$$

Taking into account that $t \tilde{u}(x) \in(0, \delta]$ and further $\delta \leq \frac{\eta_{0}}{2}$ (see (H2)(ii)), from (3.1) we have that $\theta(t \tilde{u}(x))=1$. This in addition gives

$$
\begin{equation*}
k(x, t \tilde{u})=f(x, t \tilde{u}(x))+(t \tilde{u}(x))^{p^{*}(x)-2} t \tilde{u}(x) \geq f(x, t \tilde{u}(x)) \tag{3.6}
\end{equation*}
$$

Now, since $t \tilde{u}(x) \in(0, \delta]$ and $\delta \leq \frac{\eta_{0}}{2}$, thanks to hypothesis (H2)(ii) we know that

$$
\begin{equation*}
c_{0}|t \tilde{u}|^{\tau(x)-1} \leq f(x, t \tilde{u}) . \tag{3.7}
\end{equation*}
$$

So, using (3.6) and (3.7) in (3.5) we have

$$
\begin{aligned}
& \phi_{+}(t \tilde{u}) \\
& \leq \frac{t^{p^{-}}}{p^{-}}\left[\int_{\Omega}|\nabla \tilde{u}|^{p(x)} \mathrm{d} x+\int_{\Omega} \mu(x)|\nabla \tilde{u}|^{q(x)} \mathrm{d} x\right]-\frac{c_{0} t^{\tau^{+}}}{\tau^{+}} \int_{\Omega}|\tilde{u}|^{\tau(x)} \mathrm{d} x \\
& =t^{\tau^{+}}\left[\frac{t^{p^{-}-\tau^{+}}}{p^{-}}\left[\int_{\Omega}|\nabla \tilde{u}|^{p(x)} \mathrm{d} x+\int_{\Omega} \mu(x)|\nabla \tilde{u}|^{q(x)} \mathrm{d} x\right]-\frac{c_{0}}{\tau^{+}} \int_{\Omega}|\tilde{u}|^{\tau(x)} \mathrm{d} x\right] .
\end{aligned}
$$

Now, taking into account that $\tau^{+}<p^{-}$(see (H2)(ii)), choosing $t \in(0,1)$ small enough we have that

$$
\frac{t^{p^{-}-\tau^{+}}}{p^{-}}\left[\int_{\Omega}|\nabla \tilde{u}|^{p(x)} \mathrm{d} x+\int_{\Omega} \mu(x)|\nabla \tilde{u}|^{q(x)} \mathrm{d} x\right]-\frac{c_{0}}{\tau^{+}} \int_{\Omega}|\tilde{u}|^{\tau(x)} \mathrm{d} x<0
$$

Recall that $\int_{\Omega}|\nabla \tilde{u}|^{p(x)} \mathrm{d} x+\int_{\Omega} \mu(x)|\nabla \tilde{u}|^{q(x)} \mathrm{d} x$ and $\int_{\Omega}|\tilde{u}|^{\tau(x)} \mathrm{d} x$ are fixed for a given $\tilde{u} \in \operatorname{int} C(\bar{\Omega})_{+}$. Consequently, we conclude that $\phi_{+}(t \tilde{u})<0=\phi_{+}(0)$ for $t \in(0,1)$ sufficiently small. This guarantees that $u_{0} \neq 0$.

Finally, we remark that $u_{0}$ is a global minimizer of $\phi_{+}$, so $\phi_{+}^{\prime}\left(u_{0}\right)=0$. This implies that

$$
\begin{equation*}
\int_{\Omega}\left(\left|\nabla u_{0}\right|^{p(x)-2} \nabla u_{0}+\mu(x)\left|\nabla u_{0}\right|^{q(x)-2} \nabla u_{0}\right) \cdot \nabla h \mathrm{~d} x=\int_{\Omega} k\left(x,\left(u_{0}\right)_{+}\right) h \mathrm{~d} x \tag{3.8}
\end{equation*}
$$

for all $h \in W_{0}^{1, \mathcal{H}}(\Omega)$. Now, using Proposition 2.17 of Crespo-Blanco-Gasiński-Harjulehto-Winkert [8] which gives $\pm u_{ \pm} \in W_{0}^{1, \mathcal{H}}(\Omega)$ for any $u \in W_{0}^{1, \mathcal{H}}(\Omega)$, we can choose $h=-\left(u_{0}\right)_{-}$in (3.8). In this way, we get that $\left(u_{0}\right)_{-}=0$ and thus we deduce that $u_{0} \geq 0$. As $u_{0} \neq 0$ we conclude that $u_{0}$ is a nontrivial positive weak solution of problem (3.4), thus $\mathcal{S}_{+} \neq \emptyset$. Also, from Crespo-Blanco-Winkert [9, Theorem 3.1] we have that $u_{0} \in W_{0}^{1, \mathcal{H}}(\Omega) \cap L^{\infty}(\Omega)$.

In order to obtain a nontrivial negative weak solution for problem (3.4), we can use the $C^{1}$-functional $\phi_{-}: W_{0}^{1, \mathcal{H}}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\phi_{-}(u)=\int_{\Omega}\left[\frac{1}{p(x)}|\nabla u|^{p(x)}+\frac{\mu(x)}{q(x)}|\nabla u|^{q(x)}\right] \mathrm{d} x-\int_{\Omega} K\left(x,-u_{-}\right) \mathrm{d} x
$$

for all $u \in W_{0}^{1, \mathcal{H}}(\Omega)$. Then, arguing as in the case of the positive solution, we show that it has a global minimizer which turns out to be nontrivial and nonpositive. Therefore, it is a nontrivial negative weak solution of problem (3.4).

Next, our aim is to show the existence of extremal constant sign solutions for problem (3.4). So, we consider the following auxiliary problem

$$
\begin{align*}
-\operatorname{div}\left(|\nabla u|^{p(x)-2} \nabla u+\mu(x)|\nabla u|^{q(x)-2} \nabla u\right) & =c_{0}|u|^{\tau(x)-2} u & & \text { in } \Omega,  \tag{3.9}\\
u & =0 & & \text { on } \partial \Omega
\end{align*}
$$

and establish the following result.

Proposition 3.2. Let hypothesis (H1) be satisfied. Then, problem (3.9) has a unique positive solution $\bar{u} \in W_{0}^{1, \mathcal{H}}(\Omega) \cap L^{\infty}(\Omega)$. Further, $\bar{v}=-\bar{u}$ is the unique negative solution of problem (3.9).

Proof. We point out that in order to have the claim it is sufficient to show that problem (3.9) has a unique positive solution $\bar{u} \in W_{0}^{1, \mathcal{H}}(\Omega) \cap L^{\infty}(\Omega)$. Indeed, taking into account that the problem (3.9) is odd, this guarantees that $\bar{v}=-\bar{u}$ is a negative solution of problem (3.9) and further it is the unique negative solution.

Now, in order to prove the existence of a positive solution for problem (3.9), we consider the $C^{1}$-functional $\psi_{+}: W_{0}^{1, \mathcal{H}}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\psi_{+}(u)=\int_{\Omega}\left[\frac{1}{p(x)}|\nabla u|^{p(x)}+\frac{\mu(x)}{q(x)}|\nabla u|^{q(x)}\right] \mathrm{d} x-\int_{\Omega} \frac{c_{0}}{\tau(x)}\left(u_{+}\right)^{\tau(x)} \mathrm{d} x
$$

for all $u \in W_{0}^{1, \mathcal{H}}(\Omega)$. We know that

$$
\begin{aligned}
\psi_{+}(u) & \geq \frac{1}{p^{+}} \int_{\Omega}|\nabla u|^{p(x)} \mathrm{d} x+\frac{1}{q^{+}} \int_{\Omega} \mu(x)|\nabla u|^{q(x)} \mathrm{d} x-\frac{c_{0}}{\tau^{-}} \int_{\Omega}\left(u_{+}\right)^{\tau(x)} \mathrm{d} x \\
& \geq \frac{1}{q^{+}} \rho_{\mathcal{H}}(|\nabla u|)-\frac{c_{0}}{\tau^{-}} \int_{\Omega}|u|^{\tau(x)} \mathrm{d} x
\end{aligned}
$$

With similar arguments as in the proof of Proposition 3.1, we can show that $\psi_{+}(\cdot)$ is coercive and sequentially weakly semicontinuous. This ensures, thanks to the Weierstraß-Tonelli theorem, that we can find $\bar{u} \in W_{0}^{1, \mathcal{H}}(\Omega)$ such that

$$
\psi_{+}(\bar{u})=\inf \left[\psi_{+}(u): u \in W_{0}^{1, \mathcal{H}}(\Omega)\right] .
$$

Reasoning again as in the proof of Proposition 3.1, we are also able to show that $\bar{u} \neq 0$ and further $\bar{u} \geq 0$. So, we conclude that $\bar{u} \in W_{0}^{1, \mathcal{H}}(\Omega)$ is a nontrivial positive weak solution of problem (3.9).

Next, we prove that such positive solution is unique. To this purpose, we consider the integral functional $j: L^{1}(\Omega) \rightarrow \mathbb{R} \cup\{+\infty\}$ defined by
$j(u):= \begin{cases}\int_{\Omega} \frac{1}{p(x)}\left|\nabla u^{\frac{1}{p^{-}}}\right|^{p(x)}+\int_{\Omega} \frac{\mu(x)}{q(x)}\left|\nabla u^{\frac{1}{p^{-}}}\right|^{q(x)} & \text { if } u \geq 0 \text { and } u^{\frac{1}{p^{-}}} \in W_{0}^{1, \mathcal{H}}(\Omega), \\ +\infty & \text { otherwise }\end{cases}$
and let

$$
\operatorname{dom} j=\left\{u \in L^{1}(\Omega): j(u)<+\infty\right\}
$$

be the effective domain of $j(\cdot)$. We point out that from the anisotropic Díaz-Saa inequality, see Takáč-Giacomoni [34], we have that $j(\cdot)$ is convex.

Now, let $\bar{w} \in W_{0}^{1, \mathcal{H}}(\Omega) \cap L^{\infty}(\Omega)$ be another nontrivial positive weak solution of problem (3.9). Given $\varepsilon>0$, let $\bar{u}^{\varepsilon}=\bar{u}+\varepsilon \in \operatorname{int} L^{\infty}(\Omega)_{+}$and $\bar{w}^{\varepsilon}=\bar{w}+\varepsilon \in$ int $L^{\infty}(\Omega)_{+}$. Thanks to Proposition 4.1.22 of Papageorgiou-Rădulescu-Repovš [27], we know that

$$
\begin{equation*}
\frac{\bar{u}^{\varepsilon}}{\bar{w}^{\varepsilon}} \in L^{\infty}(\Omega) \quad \text { and } \quad \frac{\bar{w}^{\varepsilon}}{\bar{u}^{\varepsilon}} \in L^{\infty}(\Omega) . \tag{3.10}
\end{equation*}
$$

We put $h=\left(\bar{u}^{\varepsilon}\right)^{p^{-}}-\left(\bar{w}^{\varepsilon}\right)^{p^{-}} \in W_{0}^{1, \mathcal{H}}(\Omega) \cap L^{\infty}(\Omega)$. We note that from (3.10) and the convexity of $j(\cdot)$ we see that the directional derivatives of $j(\cdot)$ at $\left(\bar{u}^{\varepsilon}\right)^{p^{-}}$and at
$\left(\bar{w}^{\varepsilon}\right)^{p^{-}}$in the direction $h$ exist and are equal to

$$
\begin{aligned}
j^{\prime}\left(\left(\bar{u}^{\varepsilon}\right)^{p^{-}}\right)(h) & =\frac{1}{p^{-}} \int_{\Omega} \frac{-\Delta_{p(\cdot)} \bar{u}-\mu(x) \Delta_{q(\cdot)} \bar{u}}{\left(\bar{u}^{\varepsilon}\right)^{p^{-}-1}} h \mathrm{~d} x \\
& =\frac{1}{p^{-}} \int_{\Omega} \frac{c_{0} \bar{u}^{\tau(x)-1}}{\left(\bar{u}^{\varepsilon}\right)^{p^{--1}}} h \mathrm{~d} x
\end{aligned}
$$

and

$$
\begin{aligned}
j^{\prime}\left(\left(\bar{w}^{\varepsilon}\right)^{p^{-}}\right)(h) & =\frac{1}{p^{-}} \int_{\Omega} \frac{-\Delta_{p(\cdot)} \bar{w}-\mu(x) \Delta_{q(\cdot)} \bar{w}}{\left(\bar{w}^{\varepsilon}\right)^{p^{--1}}} h \mathrm{~d} x \\
& =\frac{1}{p^{-}} \int_{\Omega} \frac{c_{0} \bar{w}^{\tau(x)-1}}{\left(\bar{w}^{\varepsilon}\right)^{p^{--1}}} h \mathrm{~d} x,
\end{aligned}
$$

respectively. Moreover, the convexity of $j(\cdot)$ implies the monotonicity of $j^{\prime}(\cdot)$. On account of this, we have that

$$
0 \leq c_{0} \int_{\Omega}\left(\frac{\bar{u}^{\tau(x)-1}}{\left(\bar{u}^{\varepsilon}\right)^{p^{-}-1}}-\frac{\bar{w}^{\tau(x)-1}}{\left(\bar{w}^{\varepsilon}\right)^{p^{--1}}}\right)\left(\left(\bar{u}^{\varepsilon}\right)^{p^{-}}-\left(\bar{w}^{\varepsilon}\right)^{p^{-}}\right) \mathrm{d} x
$$

Hence, for $\varepsilon \rightarrow 0^{+}$, using Lebesgue's Theorem, we obtain

$$
0 \leq c_{0} \int_{\Omega}\left(\frac{1}{\bar{u}^{p^{-}-\tau(x)}}-\frac{1}{\bar{w}^{p^{-}-\tau(x)}}\right)\left(\bar{u}^{p^{-}}-\bar{w}^{p^{-}}\right) \mathrm{d} x
$$

It follows that $\bar{u}=\bar{w}$ and thus we have the claim.
Finally, we point out that the double phase maximum principle leads to $\bar{u}(x)>0$ for a.a. $x \in \Omega$.

Now, we are ready to prove the existence of a smallest positive solution $u_{*} \in \mathcal{S}_{+}$ and the existence of a largest negative solution $v_{*} \in \mathcal{S}_{-}$.

Proposition 3.3. Let hypotheses (H1) and (H2) be satisfied. Then there exists $u_{*} \in \mathcal{S}_{+}$such that $u_{*} \leq u$ for all $u \in \mathcal{S}_{+}$and there exists $v_{*} \in \mathcal{S}_{-}$such that $v_{*} \geq v$ for all $v \in \mathcal{S}_{-}$.

Proof. We only show the existence of a smallest positive solution for problem (3.4), the case of a largest negative solution works similarly.

Arguing in a similar way to the proof of Proposition 7 in Papageorgiou-RădulescuRepovš [28] we can deduce that $\mathcal{S}_{+}$is downward directed. On account of this, we can use Lemma 3.10 of Hu-Papageorgiou [18, p. 178] which gives the existence of a decreasing sequence $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq \mathcal{S}_{+}$such that

$$
\inf _{n \in \mathbb{N}} u_{n}=\inf \mathcal{S}_{+}
$$

As $u_{n} \in \mathcal{S}_{+}$it follows that

$$
\begin{equation*}
\int_{\Omega}\left(\left|\nabla u_{n}\right|^{p(x)-2} \nabla u_{n}+\mu(x)\left|\nabla u_{n}\right|^{q(x)-2} \nabla u_{n}\right) \cdot \nabla h \mathrm{~d} x=\int_{\Omega} k\left(x, u_{n}\right) h \mathrm{~d} x \tag{3.11}
\end{equation*}
$$

for all $h \in W_{0}^{1, \mathcal{H}}(\Omega)$ and for all $n \in \mathbb{N}$. Hence, choosing $h=u_{n}$ in (3.11) and using (3.3) along with $0 \leq u_{n} \leq u_{1}$, we have that

$$
\rho_{\mathcal{H}}\left(\nabla u_{n}\right)=\int_{\Omega}\left|\nabla u_{n}\right|^{p(x)} \mathrm{d} x+\int_{\Omega} \mu(x)\left|\nabla u_{n}\right|^{q(x)} \mathrm{d} x \leq d_{1}
$$

for some $d_{1}>0$ and for all $n \in \mathbb{N}$. This fact along with Proposition 2.2 shows that $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W_{0}^{1, \mathcal{H}}(\Omega)$ is bounded. Hence, we can assume that

$$
u_{n} \rightharpoonup u_{*} \quad \text { in } W_{0}^{1, \mathcal{H}}(\Omega) \quad \text { and } \quad u_{n} \rightarrow u_{*} \quad \text { in } L^{q(\cdot)}(\Omega)
$$

Thanks to (3.1), (3.2) and hypothesis (H2)(ii) we have

$$
\begin{equation*}
|k(x, s)| \leq d_{2}|s|^{\tau(x)-1} \tag{3.12}
\end{equation*}
$$

for a.a. $x \in \Omega$, for all $s \in \mathbb{R}$ and for some $d_{2}>0$. From (3.12) along with $0 \leq u_{n} \leq$ $u_{1}$, we in particular deduce that

$$
\left|k\left(x, u_{n}\right)\right| \leq d_{2}\left|u_{n}\right|^{\tau(x)-1} \leq d_{2}\left|u_{1}\right|^{\tau(x)-1} .
$$

Then, taking into account that $d_{2}\left|u_{1}\right|^{\tau(\cdot)-1} \in L^{\infty}(\Omega)$, from (3.11) and (3.12) along with a Moser-iteration type argument as it was explained in Guedda-Veron [16], we can obtain that

$$
\left\|u_{n}\right\|_{\infty} \leq O\left(u_{n}\right)
$$

Now let us check that $u_{*} \neq 0$. Indeed, if $u_{*}=0$ we have that $\left\|u_{n}\right\|_{\infty} \rightarrow 0$ as $n \rightarrow+\infty$. This implies that we can find $n_{0} \in \mathbb{N}$ such that

$$
0<u_{n}(x) \leq \delta
$$

for a.a. $x \in \Omega$ and for all $n \geq n_{0}$, where $\delta \in\left(0, \min \left\{\frac{\eta_{0}}{2}, 1\right\}\right)$. Then, we fix $n \geq n_{0}$ and consider the Carathéodory function $l_{n}: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ given by

$$
l_{n}(x, s)= \begin{cases}c_{0}\left(s_{+}\right)^{\tau(x)-1} & \text { if } s \leq u_{n}(x) \\ c_{0}\left|u_{n}(x)\right|^{\tau(x)-1} & \text { if } u_{n}(x)<s\end{cases}
$$

Then we introduce the $C^{1}$-functional $\sigma_{+}: W_{0}^{1, \mathcal{H}}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\sigma_{+}(u)=\int_{\Omega}\left[\frac{1}{p(x)}|\nabla u|^{p(x)}+\frac{\mu(x)}{q(x)}|\nabla u|^{q(x)}\right] \mathrm{d} x-\int_{\Omega} L_{n}(x, u) \mathrm{d} x
$$

for all $u \in W_{0}^{1, \mathcal{H}}(\Omega)$, where $L_{n}(x, s)=\int_{0}^{s} l_{n}(x, t) \mathrm{d} t$. Similar arguments to the ones in the proof of Proposition 3.1 show that this functional is coercive and sequentially weakly lower semicontinuous. Hence, we can find $\hat{u} \in W_{0}^{1, \mathcal{H}}(\Omega)$ such that

$$
\sigma_{+}(\hat{u})=\inf \left[\sigma_{+}(u): u \in W_{0}^{1, \mathcal{H}}(\Omega)\right]<0=\sigma_{+}(0)
$$

which gives $\hat{u} \neq 0$. Also, we can deduce that $\hat{u} \geq 0$. In addition, we point out that $K_{\sigma_{+}} \subseteq\left[0, u_{n}\right]$. In fact, let $v \in K_{\sigma_{+}}$with $v \neq 0, u_{n}$, then we have

$$
\begin{equation*}
\int_{\Omega}\left(|\nabla v|^{p(x)-2} \nabla v+\mu(x)|\nabla v|^{q(x)-2} \nabla v\right) \cdot \nabla h \mathrm{~d} x=\int_{\Omega} l_{n}(x, v) h \mathrm{~d} x \tag{3.13}
\end{equation*}
$$

for all $h \in W_{0}^{1, \mathcal{H}}(\Omega)$. Taking the test function $h=\left(v-u_{n}\right)_{+}$in (3.13), using (H2)(ii) along with $0<u_{n} \leq \delta$ and $u_{n} \in S_{+}$, we obtain

$$
\begin{aligned}
& \int_{\Omega}\left(|\nabla v|^{p(x)-2} \nabla v+\mu(x)|\nabla v|^{q(x)-2} \nabla v\right) \cdot \nabla\left(v-u_{n}\right)_{+} \mathrm{d} x \\
& =\int_{\Omega} l_{n}(x, v)\left(v-u_{n}\right)_{+} \mathrm{d} x \\
& =\int_{\Omega} c_{0}\left(u_{n}\right)^{\tau(x)-1}\left(v-u_{n}\right)_{+} \mathrm{d} x
\end{aligned}
$$

$$
\begin{aligned}
& \leq \int_{\Omega} f\left(x, u_{n}\right)\left(v-u_{n}\right)_{+} \mathrm{d} x \\
& \leq \int_{\Omega}\left[f\left(x, u_{n}\right)+\left(u_{n}\right)^{p^{*}(x)-2} u_{n}\right]\left(v-u_{n}\right)_{+} \mathrm{d} x \\
& =\int_{\Omega} k\left(x, u_{n}\right)\left(v-u_{n}\right)_{+} \mathrm{d} x \\
& =\int_{\Omega}\left(\left|\nabla u_{n}\right|^{p(x)-2} \nabla u_{n}+\mu(x)\left|\nabla u_{n}\right|^{q(x)-2} \nabla u_{n}\right) \cdot \nabla\left(v-u_{n}\right)_{+} \mathrm{d} x
\end{aligned}
$$

as $0<u_{n} \leq \delta \leq \frac{\eta_{0}}{2}$ and (3.1) give $\theta\left(u_{n}\right)=1$ and so $k\left(x, u_{n}\right)=f\left(x, u_{n}\right)+$ $\left(u_{n}\right)^{p^{*}(x)-2} u_{n}$. Consequently, we have

$$
\begin{aligned}
& \int_{\Omega}\left(|\nabla v|^{p(x)-2} \nabla v-\left|\nabla u_{n}\right|^{p(x)-2} \nabla u_{n}\right) \cdot \nabla\left(v-u_{n}\right)_{+} \mathrm{d} x \\
& +\int_{\Omega} \mu(x)\left(|\nabla v|^{q(x)-2} \nabla v-\left|\nabla u_{n}\right|^{q(x)-2} \nabla u_{n}\right) \cdot \nabla\left(v-u_{n}\right)_{+} \mathrm{d} x \leq 0
\end{aligned}
$$

which implies that $v \leq u_{n}$. Furthermore, choosing the test function $h=-\left(v_{-}\right)$in (3.13), we easily see that $v_{-}=0$ and thus we have that $v \geq 0$. We finally conclude that $K_{\sigma_{+}} \subseteq\left[0, u_{n}\right]$.

Now, we recall that $\hat{u}$ is a global minimizer of $\sigma_{+}$and then $\hat{u} \in K_{\sigma_{+}} \subseteq\left[0, u_{n}\right]$. So, we know that

$$
\int_{\Omega}\left(|\nabla \hat{u}|^{p(x)-2} \nabla \hat{u}+\mu(x)|\nabla \hat{u}|^{q(x)-2} \nabla \hat{u}\right) \cdot \nabla h \mathrm{~d} x=\int_{\Omega} c_{0}\left(\hat{u}_{+}\right)^{\tau(x)-1} h \mathrm{~d} x
$$

for all $h \in W_{0}^{1, \mathcal{H}}(\Omega)$. This clearly implies that $\hat{u}$ is a positive solution of problem (3.9). Since such solution is unique (see Proposition 3.2), we conclude that $\hat{u}=\bar{u}$. From this, it follows that $\bar{u} \leq u_{n}$ for all $n \geq n_{0}$, which contradicts the hypothesis $u_{*}=0$. Consequently, we have that $u_{*} \neq 0$ and further $u_{*} \in S_{+}$. Therefore, we have that $u_{*}$ is the smallest positive solution of (3.4) in $S_{+}$.

## 4. Proof of Theorem 1.1

In this section, we use the extremal constant sign solutions $u_{*}$ and $v_{*}$ given in Proposition 3.3 in order to produce a sequence of nodal (that is, sign-changing) solutions for problem (1.1). In addition, we show that such sequence converges to 0 in $W_{0}^{1, \mathcal{H}}(\Omega)$ and in $L^{\infty}(\Omega)$.

Considering truncations of $k(x, \cdot)$ at $v_{*}(x)$ and $u_{*}(x)$, we introduce the Carathéodory function $k_{*}: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ given by

$$
k_{*}(x, s):= \begin{cases}k\left(x, v_{*}(x)\right) & \text { if } s<v_{*}(x) \\ k(x, s) & \text { if } v_{*}(x) \leq s \leq u_{*}(x) \\ k\left(x, u_{*}(x)\right) & \text { if } u_{*}(x)<s\end{cases}
$$

Let $\Phi_{*}: W_{0}^{1, \mathcal{H}}(\Omega) \rightarrow \mathbb{R}$ be the $C^{1}$-functional defined by

$$
\Phi_{*}(u)=\int_{\Omega}\left[\frac{1}{p(x)}|\nabla u|^{p(x)}+\frac{\mu(x)}{q(x)}|\nabla u|^{q(x)}\right] \mathrm{d} x-\int_{\Omega} K_{*}(x, u) \mathrm{d} x
$$

for all $u \in W_{0}^{1, \mathcal{H}}(\Omega)$, where $K_{*}(x, s)=\int_{0}^{s} k_{*}(x, t) \mathrm{d} t$.

Note that the set $K_{\Phi_{*}}$ of the critical points of the functional $\Phi_{*}$ is contained in the order interval $\left[v_{*}, u_{*}\right]$. In fact, for a given $u \in K_{\Phi_{*}} \backslash\left\{u_{*}, v_{*}\right\}$ we know that

$$
\begin{equation*}
\int_{\Omega}\left(|\nabla u|^{p(x)-2} \nabla u+\mu(x)|\nabla u|^{q(x)-2} \nabla u\right) \cdot \nabla h \mathrm{~d} x=\int_{\Omega} k_{*}(x, u) h \mathrm{~d} x \tag{4.1}
\end{equation*}
$$

holds for all $h \in W_{0}^{1, \mathcal{H}}(\Omega)$. Then, choosing $h=\left(u-u_{*}\right)_{+}$in (4.1), we have that

$$
\begin{aligned}
& \int_{\Omega}\left(|\nabla u|^{p(x)-2} \nabla u+\mu(x)|\nabla u|^{q(x)-2} \nabla u\right) \cdot \nabla\left(u-u_{*}\right)_{+} \mathrm{d} x \\
& =\int_{\Omega} k_{*}(x, u)\left(u-u_{*}\right)_{+} \mathrm{d} x \\
& =\int_{\Omega} k\left(x, u_{*}\right)\left(u-u_{*}\right)_{+} \mathrm{d} x \\
& =\int_{\Omega}\left(\left|\nabla u_{*}\right|^{p(x)-2} \nabla u_{*}+\mu(x)\left|\nabla u_{*}\right|^{q(x)-2} \nabla u_{*}\right) \cdot \nabla\left(u-u_{*}\right)_{+} \mathrm{d} x
\end{aligned}
$$

since $u_{*} \in \mathcal{S}_{+}$. Hence, we deduce that

$$
\begin{aligned}
& \int_{\Omega}\left(|\nabla u|^{p(x)-2} \nabla u-\left|\nabla u_{*}\right|^{p(x)-2} \nabla u_{*}\right) \cdot \nabla\left(u-u_{*}\right)_{+} \mathrm{d} x \\
& +\int_{\Omega} \mu(x)\left(|\nabla u|^{q(x)-2} \nabla u-\left|\nabla u_{*}\right|^{q(x)-2} \nabla u_{*}\right) \cdot \nabla\left(u-u_{*}\right)_{+} \mathrm{d} x=0,
\end{aligned}
$$

and consequently we have that $u \leq u_{*}$. An analogous reasoning and the choice of $h=\left(v_{*}-u\right)_{+}$in (4.1) gives that $v_{*} \leq u$.

Now, let $V \subseteq W_{0}^{1, \mathcal{H}}(\Omega) \cap L^{\infty}(\Omega)$ be a finite dimensional subspace. Then, given $v \in V$ we put

$$
\begin{aligned}
\left\{v<v_{*}\right\} & :=\left\{x \in \Omega: v(x)<v_{*}(x)\right\} \\
\left\{v_{*} \leq v \leq u_{*}\right\} & :=\left\{x \in \Omega: v_{*}(x) \leq v(x) \leq u_{*}(x)\right\} \\
\left\{u_{*}<v\right\} & :=\left\{x \in \Omega: u_{*}(x)<v(x)\right\}
\end{aligned}
$$

Since $V$ is finite dimensional, all norms on $V$ are equivalent, see for example Papageorgiou-Winkert [32, Proposition 3.1.17, p.183]. Hence, we know that there exists a positive constant $e_{V}$, independent of $v \in V$, such that

$$
\begin{equation*}
e_{V}\|v\| \leq\|v\|_{\tau(\cdot)} \tag{4.2}
\end{equation*}
$$

Using this fact, we are able to establish the following result.
Proposition 4.1. Let hypotheses (H1) and (H2) be satisfied. Then, we can find $r_{V}>0$ such that

$$
\sup \left[\Phi_{*}(v): v \in V,\|v\|=r_{V}\right]<0
$$

Proof. Since all norms on $V$ are equivalent, we can find $r_{V}>0$ such that

$$
v \in V \quad \text { and } \quad\|v\|=r_{V} \quad \text { imply } \quad|v(x)| \leq \delta \quad \text { for a.a. } x \in \Omega
$$

with $\delta \in\left(0, \min \left\{\frac{\eta_{0}}{2}, 1\right\}\right)$. Recall that from $\delta \leq \frac{\eta_{0}}{2}$ it follows $\theta(v(x))=1$ for a.a. $x \in \Omega$, see (3.1). Hence, given $v \in V$ with $\|v\|=r_{V}$ we know that

$$
k_{*}(x, v(x))= \begin{cases}f\left(x, v_{*}(x)\right)+\left|v_{*}(x)\right|^{p^{*}(x)-2} v_{*}(x) & \text { if } v(x)<v_{*}(x)  \tag{4.3}\\ f(x, v(x))+|v(x)|^{p^{*}(x)-2} v(x) & \text { if } v_{*}(x) \leq v(x) \leq u_{*}(x) \\ f\left(x, u_{*}(x)\right)+\left|u_{*}(x)\right|^{p^{*}(x)-2} u_{*}(x) & \text { if } u_{*}(x)<v(x)\end{cases}
$$

Let $f_{*}: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ be the function defined by

$$
f_{*}(x, v(x))= \begin{cases}f\left(x, v_{*}(x)\right) & \text { if } v(x)<v_{*}(x) \\ f(x, v(x)) & \text { if } v_{*}(x) \leq v(x) \leq u_{*}(x) \\ f\left(x, u_{*}(x)\right) & \text { if } u_{*}(x)<v(x)\end{cases}
$$

and in addition let $F_{*}(x, s):=\int_{0}^{s} f_{*}(x, t) \mathrm{d} t$. Then, for $v<v_{*}$ we see that

$$
\begin{aligned}
F_{*}(x, v) & =\int_{0}^{v_{*}} f_{*}(x, s) \mathrm{d} s+\int_{v_{*}}^{v} f_{*}(x, s) \mathrm{d} s \\
& =\int_{0}^{v_{*}} f(x, s) \mathrm{d} s+f\left(x, v_{*}\right)\left(v-v_{*}\right)
\end{aligned}
$$

Note that according to (H2)(ii) we have

$$
\begin{equation*}
f(x, s)<0 \text { for all }-\delta<s \leq 0 \quad \text { and } \quad f(x, s)>0 \text { for all } 0<s \leq \delta \tag{4.4}
\end{equation*}
$$

Using (4.4), we see that $f\left(x, v_{*}\right)$ is negative and hence we infer that $f\left(x, v_{*}\right)(v-$ $\left.v_{*}\right)>0$. Therefore, we can write

$$
\begin{aligned}
F(x, v)-F_{*}(x, v) & =F(x, v)-F\left(x, v_{*}\right)+f\left(x, v_{*}\right)\left(v_{*}-v\right) \\
& \leq F(x, v)-F\left(x, v_{*}\right),
\end{aligned}
$$

where $F(x, s):=\int_{0}^{s} f(x, t) \mathrm{d} t$. Moreover, for $u_{*}<v$ we have that

$$
F_{*}(x, v)=F\left(x, u_{*}\right)+f\left(x, u_{*}\right)\left(v-u_{*}\right) .
$$

Taking into account that from (4.4) it follows $f\left(x, u_{*}\right)\left(u_{*}-v\right)<0$, we deduce that

$$
\begin{aligned}
F(x, v)-F_{*}(x, v) & =F(x, v)-F\left(x, u_{*}\right)+f\left(x, u_{*}\right)\left(u_{*}-v\right) \\
& \leq F(x, v)-F\left(x, u_{*}\right) .
\end{aligned}
$$

Using this facts along with the fact that the terms

$$
\left.\frac{1}{p^{*}(x)}\left|v_{*}\right|\right|^{p^{*}(x)}, \quad \frac{1}{p^{*}(x)}|v|^{p^{*}(x)} \quad \text { and }\left.\quad \frac{1}{p^{*}(x)}\left|u_{*}\right|\right|^{p^{*}(x)}
$$

are positive, we get that

$$
\begin{aligned}
\Phi_{*}(v)= & \int_{\Omega}\left[\frac{1}{p(x)}|\nabla v|^{p(x)}+\frac{\mu(x)}{q(x)}|\nabla v|^{q(x)}\right] \mathrm{d} x-\int_{\Omega} K_{*}(x, v) \mathrm{d} x \\
\leq & \frac{1}{p^{-}} \int_{\Omega}|\nabla v|^{p(x)} \mathrm{d} x+\frac{1}{q^{-}} \int_{\Omega} \mu(x)|\nabla v|^{q(x)} \mathrm{d} x \\
& -\int_{\left\{v<v_{*}\right\}}\left[F_{*}(x, v)+\frac{1}{p^{*}(x)}\left|v_{*}\right|^{p^{*}(x)}\right] \mathrm{d} x \\
& -\int_{\left\{v_{*} \leq v \leq u_{*}\right\}}\left[F(x, v)+\frac{1}{p^{*}(x)}|v|^{p^{*}(x)}\right] \mathrm{d} x \\
& -\int_{\left\{u_{*}<v\right\}}\left[F_{*}(x, v)+\frac{1}{p^{*}(x)}\left|u_{*}\right|^{p^{*}(x)}\right] \mathrm{d} x \\
\leq & \frac{1}{p^{-}} \int_{\Omega}|\nabla v|^{p(x)} \mathrm{d} x+\frac{1}{q^{-}} \int_{\Omega} \mu(x)|\nabla v|^{q(x)} \mathrm{d} x \\
& -\int_{\left\{v<v_{*}\right\}} F_{*}(x, v) \mathrm{d} x-\int_{\left\{v_{*} \leq v \leq u_{*}\right\}} F(x, v) \mathrm{d} x-\int_{\left\{u_{*}<v\right\}} F_{*}(x, v) \mathrm{d} x
\end{aligned}
$$

$$
\begin{aligned}
\leq & \frac{1}{p^{-}} \int_{\Omega}|\nabla v|^{p(x)} \mathrm{d} x+\frac{1}{q^{-}} \int_{\Omega} \mu(x)|\nabla v|^{q(x)} \mathrm{d} x-\int_{\Omega} F(x, v) \mathrm{d} x \\
& +\int_{\left\{v<v_{*}\right\}}\left[F(x, v)-F_{*}(x, v)\right] \mathrm{d} x+\int_{\left\{u_{*}<v\right\}}\left[F(x, v)-F_{*}(x, v)\right] \mathrm{d} x \\
\leq & \frac{1}{p^{-}} \int_{\Omega}|\nabla v|^{p(x)} \mathrm{d} x+\frac{1}{q^{-}} \int_{\Omega} \mu(x)|\nabla v|^{q(x)} \mathrm{d} x-\int_{\Omega} F(x, v) \mathrm{d} x \\
& +\int_{\left\{v<v_{*}\right\}}\left[F(x, v)-F\left(x, v_{*}\right)\right] \mathrm{d} x+\int_{\left\{u_{*}<v\right\}}\left[F(x, v)-F\left(x, u_{*}\right)\right] \mathrm{d} x \\
\leq & \frac{1}{p^{-}} \rho_{\mathcal{H}}(\nabla v)-\int_{\Omega} F(x, v) \mathrm{d} x \\
& +\int_{\left\{v<v_{*}\right\}}\left[F(x, v)-F\left(x, v_{*}\right)\right] \mathrm{d} x+\int_{\left\{u_{*}<v\right\}}\left[F(x, v)-F\left(x, u_{*}\right)\right] \mathrm{d} x .
\end{aligned}
$$

Now, due to hypothesis (H2)(ii) we can find $\delta \in\left(0, \min \left\{\frac{\eta_{0}}{2}, 1\right\}\right)$ such that

$$
c_{0}|s|^{\tau(x)-1} \leq|f(x, s)| \quad \text { for all }|s| \leq \delta
$$

Consequently, as $\tau^{+}<p^{-}$due to hypothesis (H2)(ii), we can choose

$$
r_{V}<\left(e_{V}\right)^{\frac{\tau^{+}}{p^{-}-\tau^{+}}}
$$

with $e_{V} r_{V}<1$, where $e_{V}$ is the positive constant introduced in (4.2), so that

$$
\int_{\left\{v<v_{*}\right\}}\left[F(x, v)-F\left(x, v_{*}\right)\right] \mathrm{d} x+\int_{\left\{u_{*}<v\right\}}\left[F(x, v)-F\left(x, u_{*}\right)\right] \mathrm{d} x<\frac{c_{0}\left(e_{V} r_{V}\right)^{\tau^{+}}}{2 \tau^{+}}
$$

and in addition

$$
\Phi_{*}(v) \leq \frac{1}{p^{-}} \rho_{\mathcal{H}}(\nabla v)-\frac{c_{0}}{\tau^{+}} \int_{\Omega}|v|^{\tau(x)} \mathrm{d} x+\frac{c_{0}\left(e_{V} r_{V}\right)^{\tau^{+}}}{2 \tau^{+}}
$$

Next, using Proposition 2.2 (ii), (iii) we have

$$
\rho_{\mathcal{H}}(\nabla v) \leq \max \left\{\|v\|^{p^{-}},\|v\|^{q^{+}}\right\}
$$

and using Proposition 2.1 (ii), (iii) we have

$$
\rho_{\tau(\cdot)}(v) \geq \min \left\{\|v\|_{\tau(x)}^{\tau^{-}},\|v\|_{\tau(x)}^{\tau^{+}}\right\}
$$

Further, thanks to (4.2) we deduce that

$$
\rho_{\tau(\cdot)}(v) \geq \min \left\{\left(e_{V}\|v\|\right)^{\tau^{-}},\left(e_{V}\|v\|\right)^{\tau^{+}}\right\}
$$

Now, as $r_{V}<1$ and $e_{V} r_{V}<1$, for $v \in V$ with $\|v\|=r_{V}$, we have

$$
\rho_{\mathcal{H}}(\nabla v) \leq r_{V}^{p^{-}} \quad \text { and } \quad \rho_{\tau(\cdot)}(v) \geq\left(e_{V} r_{V}\right)^{\tau^{+}} .
$$

Hence, we get

$$
\begin{aligned}
\Phi_{*}(v) & \leq \frac{1}{p^{-}} r_{V}^{p^{-}}-\frac{c_{0}\left(e_{V} r_{V}\right)^{\tau^{+}}}{\tau^{+}}+\frac{c_{0}\left(e_{V} r_{V}\right)^{\tau^{+}}}{2 \tau^{+}} \\
& =\left(\frac{r_{V}^{p^{-}-\tau^{+}}}{p^{-}}-\frac{c_{0} e_{V}^{\tau^{+}}}{2 \tau^{+}}\right) r_{V}^{\tau^{+}}
\end{aligned}
$$

Now, we recall that $c_{0}>\frac{2 \tau^{+}}{p^{-}}($see $(H 2)(i i))$ and $r_{V}<\left(e_{V}\right)^{\frac{\tau^{+}}{p^{-}-\tau^{+}}}$. Therefore, we have that

$$
\frac{r_{V}^{p^{-}-\tau^{+}}}{p^{-}}-\frac{c_{0} e_{V}^{\tau^{+}}}{2 \tau^{+}}<\frac{r_{V}^{p^{-}-\tau^{+}}}{p^{-}}-\frac{e_{V}^{\tau^{+}}}{p^{-}}<0 .
$$

From this, we conclude that $\Phi_{*}(v)<0$ for all $v \in V$ with $\|v\|=r_{V}$ and $r_{V}<$ $\left(e_{V}\right)^{\frac{\tau^{+}}{p^{-}-\tau^{+}}}$such that $e_{V} r_{V}<1$. This clearly gives the assertion of the proposition.

We are now ready to establish the proof of Theorem 1.1 which is based on a generalized version of the symmetric mountain pass theorem due to Kajikiya [19, Theorem 1].

Proof of Theorem 1.1. According to the definition of $k_{*}: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ given in (4.3), we can show with similar arguments as in the proof of Proposition 3.1 that the functional $\Phi_{*}$ is even and coercive. This permits us to deduce that $\Phi_{*}$ is bounded from below. Further, from Papageorgiou-Rădulescu-Repovš [27, Proposition 5.1.15], we also have that $\Phi_{*}$ satisfies the PS-condition. Therefore, taking into account that Proposition 4.1 holds, we can use Theorem 1 of Kajikiya [19] which guarantees the existence of a sequence $\left\{z_{n}\right\}_{n \in \mathbb{N}} \subset W_{0}^{1, \mathcal{H}}(\Omega) \cap L^{\infty}(\Omega)$ satisfying the following properties

$$
z_{n} \in K_{\Phi_{*}} \subseteq\left[v_{*}, u_{*}\right], \quad z_{n} \neq 0, \quad \Phi_{*}\left(z_{n}\right) \leq 0 \quad \text { for all } n \in \mathbb{N}
$$

and

$$
\left\|z_{n}\right\| \rightarrow 0 \quad \text { as } n \rightarrow+\infty
$$

Since $v_{*}$ and $u_{*}$ are extremal solutions of problem (3.4), from

$$
z_{n} \in K_{\Phi_{*}} \subseteq\left[v_{*}, u_{*}\right] \quad \text { and } \quad z_{n} \neq 0 \quad \text { for all } n \in \mathbb{N}
$$

we conclude that $z_{n}$ is a nodal (that is, sign-changing) solution of problem (3.4) for all $n \in \mathbb{N}$. Now, we recall that

$$
\left\|z_{n}\right\|_{\infty} \leq O\left(u_{n}\right)
$$

(see the proof of Proposition 3.3). This along with $\left\|z_{n}\right\| \rightarrow 0$ yields $\left\|z_{n}\right\|_{\infty} \rightarrow 0$ as $n \rightarrow+\infty$.

Finally, from $\left\|z_{n}\right\|_{\infty} \rightarrow 0$ as $n \rightarrow+\infty$ we also get that there exists $n_{0} \in \mathbb{N}$ such that $\left|z_{n}(x)\right| \leq \frac{\eta_{0}}{2}$ for a.a. $x \in \Omega$ and for all $n \geq n_{0}$. Thus, we have that $\theta\left(z_{n}(x)\right)=1$ for a.a. $x \in \Omega$ and for all $n \geq n_{0}$. On account of this, in view of (3.2), we conclude that $z_{n}$ is a nodal solution of problem (1.1) for all $n \geq n_{0}$.

## References

[1] A. Aberqi, J. Bennouna, O. Benslimane, M.A. Ragusa, Existence results for double phase problem in Sobolev-Orlicz spaces with variable exponents in complete manifold, Mediterr. J. Math. 19 (2022), no. 4, Paper No. 158, 19 pp.
[2] K.S. Albalawi, N.H. Alharthi, F. Vetro, Gradient and parameter dependent Dirichlet $(p(x), q(x))$-Laplace type problem, Mathematics 10 (2022), no. 8, Article number 1336.
[3] A. Bahrouni, V.D. Rădulescu, P. Winkert, Double phase problems with variable growth and convection for the Baouendi-Grushin operator, Z. Angew. Math. Phys. 71 (2020), no. 6, 183, 14 pp.
[4] P. Baroni, M. Colombo, G. Mingione, Harnack inequalities for double phase functionals, Nonlinear Anal. 121 (2015), 206-222.
[5] P. Baroni, M. Colombo, G. Mingione, Non-autonomous functionals, borderline cases and related function classes, St. Petersburg Math. J. 27 (2016), 347-379.
[6] P. Baroni, M. Colombo, G. Mingione, Regularity for general functionals with double phase, Calc. Var. Partial Differential Equations 57 (2018), no. 2, Art. 62, 48 pp.
[7] F. Colasuonno, M. Squassina, Eigenvalues for double phase variational integrals, Ann. Mat. Pura Appl. (4) 195 (2016), no. 6, 1917-1959.
[8] Á. Crespo-Blanco, L. Gasiński, P. Harjulehto, P. Winkert, A new class of double phase variable exponent problems: Existence and uniqueness, J. Differential Equations 323 (2022), 182-228.
[9] Á. Crespo-Blanco, P. Winkert, Nehari manifold approach for superlinear double phase problems with variable exponents, Ann. Mat. Pura Appl. (4), https://doi.org/10.1007/s10231-023-01375-2.
[10] C. De Filippis, G. Mingione, Lipschitz bounds and nonautonomous integrals, Arch. Ration. Mech. Anal. 242 (2021), 973-1057.
[11] L. Diening, P. Harjulehto, P. Hästö, M. Růžička, Lebesgue and Sobolev Spaces with Variable Exponents, Springer, Heidelberg, 2011.
[12] C. Farkas, P. Winkert, An existence result for singular Finsler double phase problems, J. Differential Equations 286 (2021), 455-473.
[13] L. Gasiński, N.S. Papageorgiou, Constant sign and nodal solutions for superlinear double phase problems, Adv. Calc. Var. 14 (2021), no. 4, 613-626.
[14] L. Gasiński, P. Winkert, Constant sign solutions for double phase problems with superlinear nonlinearity, Nonlinear Anal. 195 (2020), 111739.
[15] L. Gasiński, P. Winkert, Existence and uniqueness results for double phase problems with convection term, J. Differential Equations 268 (2020), no. 8, 4183-4193.
[16] M. Guedda, L. Véron, Quasilinear elliptic equations involving critical Sobolev exponents, Nonlinear Anal. 13 (1989), no. 8, 879-902.
[17] P. Harjulehto, P. Hästö, "Orlicz Spaces and Generalized Orlicz Spaces", Springer, Cham, 2019.
[18] S. Hu, N.S. Papageorgiou, "Handbook of Multivalued Analysis. Vol. I", Kluwer Academic Publishers, Dordrecht, 1997.
[19] R. Kajikiya, A critical point theorem related to the symmetric mountain pass lemma and its applications to elliptic equations, J. Funct. Anal. 225 (2005), no. 2, 352-370.
[20] I.H. Kim, Y.-H. Kim, M.W. Oh, S. Zeng, Existence and multiplicity of solutions to concave-convex-type double-phase problems with variable exponent, Nonlinear Anal. Real World Appl. 67 (2022), Paper No. 103627, 25 pp.
[21] S. Leonardi, N.S. Papageorgiou, Anisotropic Dirichlet double phase problems with competing nonlinearities, Rev. Mat. Complut. 36 (2023), no. 2, 469-490.
[22] W. Liu, G. Dai, Existence and multiplicity results for double phase problem, J. Differential Equations 265 (2018), no. 9, 4311-4334.
[23] Z. Liu, N.S. Papageorgiou, Asymptotically vanishing nodal solutions for critical double phase problems, Asymptot. Anal. 124 (2021), no. 3-4, 291-302.
[24] J. Liu, P. Pucci, Existence of solutions for a double-phase variable exponent equation without the Ambrosetti-Rabinowitz condition, Adv. Nonlinear Anal. 12 (2023), no. 1, Paper No. 20220292, 18 pp.
[25] P. Marcellini, Regularity of minimizers of integrals of the calculus of variations with nonstandard growth conditions, Arch. Rational Mech. Anal. 105 (1989), no. 3, 267-284.
[26] P. Marcellini, Regularity and existence of solutions of elliptic equations with $p, q$-growth conditions, J. Differential Equations 90 (1991), no. 1, 1-30.
[27] N.S. Papageorgiou, V.D. Rădulescu, D.D. Repovš, "Nonlinear Analysis-Theory and Methods", Springer, Cham, 2019.
[28] N.S. Papageorgiou, V.D. Rădulescu, D.D. Repovš, Positive solutions for perturbations of the Robin eigenvalue problem plus an indefinite potential, Discrete Contin. Dyn. Syst. 37 (2017), no. 5, 2589-2618.
[29] N.S. Papageorgiou, C. Vetro, Superlinear $(p(z), q(z))$-equations, Complex Var. Elliptic Equ. 64 (2019), no. 1, 8-25.
[30] N.S. Papageorgiou, C. Vetro, F. Vetro, Solutions for parametric double phase Robin problems, Asymptot. Anal. 121 (2021), no. 2, 159-170.
[31] N.S. Papageorgiou, F. Vetro, P. Winkert, Sign changing solutions for critical double phase problems with variable exponent, Z. Anal. Anwend. 42 (2023), no. 1-2, 235-251.
[32] N.S. Papageorgiou, P. Winkert, "Applied Nonlinear Functional Analysis", De Gruyter, Berlin, 2018.
[33] K. Perera, M. Squassina, Existence results for double-phase problems via Morse theory, Commun. Contemp. Math. 20 (2018), no. 2, 1750023, 14 pp.
[34] P. Takáč, J. Giacomoni, A $p(x)$-Laplacian extension of the Díaz-Saa inequality and some applications, Proc. Roy. Soc. Edinburgh Sect. A 150 (2020), no. 1, 205-232.
[35] F. Vetro, P. Winkert, Constant sign solutions for double phase problems with variable exponents, Appl. Math. Lett. 135 (2023), Paper No. 108404, 7 pp.
[36] F. Vetro, P. Winkert, Existence, uniqueness and asymptotic behavior of parametric anisotropic ( $p, q$ )-equations with convection, Appl. Math. Optim. 86 (2022), no. 2, Paper No. 18, 18 pp.
[37] F. Vetro, P. Winkert, Nodal solutions for critical Robin double phase problems with variable exponent, Discrete Contin. Dyn. Syst. Ser. S 16 (2023), no. 11, 3333-3349.
[38] S. Zeng, Y. Bai, L. Gasiński, P. Winkert, Existence results for double phase implicit obstacle problems involving multivalued operators, Calc. Var. Partial Differential Equations 59 (2020), no. 5,176 .
[39] S. Zeng, V.D. Rădulescu, P. Winkert, Double phase obstacle problems with variable exponent, Adv. Differential Equations 27 (2022), no. 9-10, 611-645.
[40] V.V. Zhikov, Averaging of functionals of the calculus of variations and elasticity theory, Izv. Akad. Nauk SSSR Ser. Mat. 50 (1986), no. 4, 675-710.
[41] V.V. Zhikov, On Lavrentiev's phenomenon, Russian J. Math. Phys. 3 (1995), no. 2, 249-269.
[42] V.V. Zhikov, On variational problems and nonlinear elliptic equations with nonstandard growth conditions, J. Math. Sci. 173 (2011), no. 5, 463-570.
(N.S. Papageorgiou) National Technical University, Department of Mathematics, Zografou Campus, Athens 15780, Greece

Email address: npapg@math.ntua.gr
(F. Vetro) Palermo, 90123 Palermo, Italy

Email address: francescavetro80@gmail.com
(P. Winkert) Technische Universität Berlin, Institut für Mathematik, Strasse des 17. Juni 136, 10623 Berlin, Germany

Email address: winkert@math.tu-berlin.de


[^0]:    2020 Mathematics Subject Classification. 35A01, 35D30, 35J60, 35J62, 35J66.
    Key words and phrases. Critical problem, double phase operator, existence results, multiple solutions, nodal solutions, sign-changing solutions, variable exponent.

